

You may drop one 12-point problem, but you have to solve *all* problems worth more than 12 points.

1. (12pts) Find all equilibria, and categorize their stability. Make two plots: (1) phase plot (dy/dt vs. y); (2) plot the solution for the given initial condition, $y(t)$ vs t .

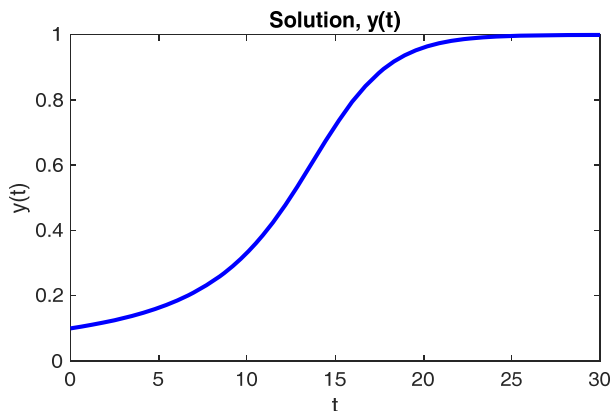
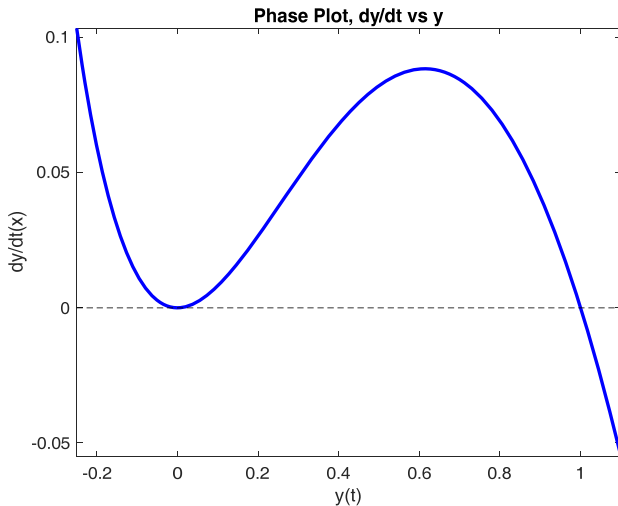
$$\begin{cases} \frac{dy}{dt} = (1-y)[\ln(1+y)]^2 \\ y(0) = 0.1 \end{cases}$$

1. $y_{eq} = 0$ $\ln(1+y) = y + O(y^2) \Rightarrow \frac{dy}{dt} \approx (1-y)y^2 \Rightarrow \frac{dy}{dt} \approx y^2 \Rightarrow y_{eq} = 0$ is non-hyperbolic, semi-stable

2. $y_{eq} = 1$ denote $Y = y - 1 \Rightarrow \frac{dY}{dt} \approx \alpha Y$ where $\alpha = -\ln^2 2 < 0 \Rightarrow$ linearly (therefore, asymptotically) stable

Note: you could use direct differentiation to prove that $f'(1) = -\ln^2 2$, but that's unnecessary, and causes mistakes

Instead, note that near $y = 1$ we have $\ln(1+y) = \ln 2 + O(y-1)$



2. (12pts) Consider the so-called RLC electric circuit equation (you don't have to know what it means):

$$L \frac{d^2 q}{dt^2} + R \frac{dq}{dt} + \frac{q(t)}{C} = \Phi(t)$$

Fundamental units are: $[t] = T$ (time), $[q] = Q$ (charge), $[\Phi] = V$ (electric potential)

- Determine the fundamental units of constants R , L and C (resistance, inductance and capacitance).
- Find any two distinct time scales t_c , in terms of model parameters. You don't have to use linear algebra.
- Explain why you can only eliminate two parameters by non-dimensionalization in this case, not three.
- Non-dimensionalize this equation, using electron charge e as an extra scale: $q_c = e$, $\bar{q}(t) = \frac{q(t)}{e}$.

$$L \frac{d^2 q}{dt^2} + R \frac{dq}{dt} + \frac{q(t)}{C} = \Phi(t) \Rightarrow [L] \frac{Q}{T^2} + [R] \frac{Q}{T} + \frac{Q}{[C]} = V \Rightarrow \begin{cases} [L] = \frac{V}{Q} T^2 \\ [R] = \frac{V}{Q} T \\ [C] = \frac{Q}{V} \end{cases} \Rightarrow \begin{cases} t_c = RC \\ t_c = \sqrt{LC} \\ t_c = \frac{L}{R} \end{cases} \quad (2 \text{ out of 3 is enough})$$

- Can eliminate 2 parameters, since only $N_U = 2$ units are independent, $\frac{V}{Q}$ and T : $N_D - N_U = 6 - 2 = N_{\Pi} = 4$
- Of course, this is only true in the context of this problem: Q and V are, in general, independent units
- Start by non-dimensionalizing the charge, the potential, and the entire equation:

$$L \frac{d^2 q}{dt^2} + R \frac{dq}{dt} + \frac{q(t)}{C} = \Phi(t) \quad \left\| \times \frac{C}{e} \Rightarrow LC \frac{d^2 \bar{q}}{dt^2} + RC \frac{d\bar{q}}{dt} + \bar{q} = \frac{C\Phi}{e} = \bar{\Phi}$$

- Three simplest non-dimensionalization options, depending on time scale choices (either of these is fine):

$$1. t_c = RC \Rightarrow t = RC \bar{t} \Rightarrow \frac{LC}{(RC)^2} \frac{d^2 \bar{q}}{d\bar{t}^2} + \frac{d\bar{q}}{d\bar{t}} + \bar{q} = \bar{\Phi} \Rightarrow \boxed{\rho \frac{d^2 \bar{q}}{d\bar{t}^2} + \frac{d\bar{q}}{d\bar{t}} + \bar{q} = \bar{\Phi}} \quad \text{where } \rho = \frac{L}{R^2 C}$$

$$2. t_c = \sqrt{LC} \Rightarrow t = \sqrt{LC} \bar{t} \Rightarrow \frac{d^2 \bar{q}}{d\bar{t}^2} + \frac{RC}{\sqrt{LC}} \frac{d\bar{q}}{d\bar{t}} + \bar{q} = \bar{\Phi} \Rightarrow \boxed{\frac{d^2 \bar{q}}{d\bar{t}^2} + \rho \frac{d\bar{q}}{d\bar{t}} + \bar{q} = \bar{\Phi}} \quad \text{where } \rho = R \sqrt{\frac{C}{L}}$$

$$3. t_c = \frac{L}{R} \Rightarrow t = \frac{L}{R} \bar{t} \Rightarrow LC \frac{R^2}{L^2} \frac{d^2 \bar{q}}{d\bar{t}^2} + RC \frac{R}{L} \frac{d\bar{q}}{d\bar{t}} + \bar{q} = \bar{\Phi} \quad \left\| \times \frac{L}{CR^2} \Rightarrow \boxed{\frac{d^2 \bar{q}}{d\bar{t}^2} + \frac{d\bar{q}}{d\bar{t}} + \rho \bar{q} = \bar{\Phi}} \quad \text{where } \rho = \frac{L}{R^2 C}; \quad \bar{\Phi} = \frac{C\Phi}{e} \frac{L}{CR^2} = \frac{L\Phi}{R^2 e}$$

3. (12pts) Consider a charged ball of radius R with spherically-symmetric charge density equal to

$$\rho(\mathbf{r}) = \rho(r) = \begin{cases} \frac{\gamma}{r}, & r \leq R \quad (\gamma = \text{const}) \\ 0, & r > R \end{cases}$$

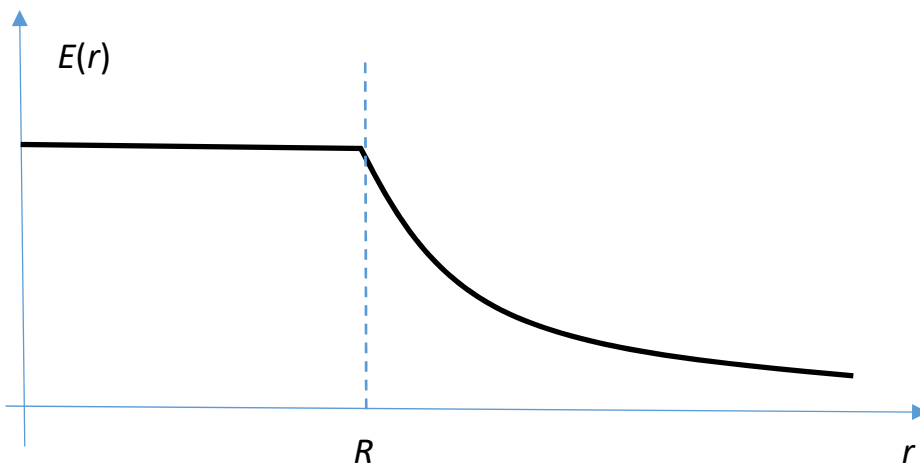
Apply the divergence theorem to the Gauss law $\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$ to find the electric field $\mathbf{E}(r)$ both inside and outside of this ball. Plot $E(r) = |\mathbf{E}|$ as a function of r . Make sure to explain all steps clearly

$$\iiint_{\text{BALL}} \nabla \cdot \mathbf{E}(\mathbf{r}) dV = \iint_{\text{SPHERE}} \mathbf{E}(\mathbf{r}) \cdot \mathbf{n} dS = \iint_{\text{SPHERE}} E(r) (\hat{\mathbf{r}} \cdot \hat{\mathbf{r}}) dS = E(r) \cdot 4\pi r^2$$

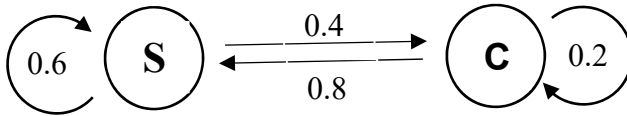
$$= \frac{1}{\epsilon_0} \iiint_{\text{BALL}} \rho(\mathbf{r}) dV = \begin{cases} r \leq R: & = \frac{1}{\epsilon_0} 4\pi \int_0^r \frac{\gamma}{\hat{r}} \hat{r}^2 d\hat{r} = \frac{4\pi\gamma}{\epsilon_0} \int_0^r \hat{r} d\hat{r} = \frac{4\pi\gamma r^2}{2\epsilon_0} \\ r > R: & \frac{1}{\epsilon_0} 4\pi \int_0^R \frac{\gamma}{\hat{r}} \hat{r}^2 d\hat{r} = \frac{4\pi\gamma}{\epsilon_0} \int_0^R \hat{r} d\hat{r} = \frac{4\pi\gamma R^2}{2\epsilon_0} \end{cases}$$

$$\Rightarrow E(r) = \begin{cases} r \leq R: & \frac{4\pi\gamma r^2}{2\epsilon_0} \frac{1}{4\pi r^2} = \frac{\gamma}{2\epsilon_0} = \text{const} \\ r > R: & \frac{4\pi\gamma R^2}{2\epsilon_0} \frac{1}{4\pi r^2} = \frac{\gamma}{2\epsilon_0} \frac{R^2}{r^2} \end{cases}$$

Solution is finite, despite unbounded density



4. (12pts) Consider the **discrete state, discrete time** Markov Chain describing a toy weather model, with daily transitions between “S” (sunny) and “C” (cloudy) days (supposedly obtained using repeated observation):



- a) Write down the explicit solution of this discrete-time dynamical system, assuming that the weather was cloudy on day zero
- b) What is the probability that it is sunny on day 4, given that it is cloudy on day zero? One decimal digit of precision is enough in your answer.

$$M = \begin{pmatrix} 0.6 & 0.8 \\ 0.4 & 0.2 \end{pmatrix} \begin{matrix} \text{S} \\ \text{C} \end{matrix} \quad \text{trace}(M) = 0.8 = \lambda_1 + \lambda_2 = 1 + \lambda_2 \Rightarrow \lambda_2 = 0.8 - 1 = -0.2$$

$$\boxed{\lambda_1 = 1} \Rightarrow (M - \lambda_1 I) \mathbf{v}_1 = \begin{pmatrix} -0.4 & 0.8 \\ 0.4 & -0.8 \end{pmatrix} \mathbf{v}_1 = 0 \Rightarrow \mathbf{v}_1 = \begin{pmatrix} 2 \\ 1 \end{pmatrix} \Rightarrow \mathbf{p}^{EQ} = c_1 \mathbf{v}_1 = \begin{pmatrix} 2/3 \\ 1/3 \end{pmatrix}$$

$$\boxed{\lambda_2 = -0.2} \Rightarrow (M - \lambda_2 I) \mathbf{v}_2 = \begin{pmatrix} 0.8 & 0.8 \\ 0.4 & 0.4 \end{pmatrix} \mathbf{v}_2 = 0 \Rightarrow \mathbf{v}_2 = \begin{pmatrix} 1 \\ -1 \end{pmatrix}$$

$$\text{Use initial condition: } \mathbf{p}^0 = c_1 \mathbf{v}_1 + c_2 \mathbf{v}_2 = \mathbf{p}^{EQ} + c_2 \mathbf{v}_2 = \begin{pmatrix} 2/3 \\ 1/3 \end{pmatrix} + c_2 \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \Rightarrow \boxed{c_2 = -\frac{2}{3}}$$

$$\text{Therefore, the solution is } \mathbf{p}^m = c_1 \lambda_1^m \mathbf{v}_1 + c_2 \lambda_2^m \mathbf{v}_2 = \begin{pmatrix} 2/3 \\ 1/3 \end{pmatrix} - \frac{2}{3} (-0.2)^m \begin{pmatrix} 1 \\ -1 \end{pmatrix} = \begin{pmatrix} \frac{2}{3} [1 - (-0.2)^m] \\ \frac{1}{3} [1 - 2(-0.2)^m] \end{pmatrix}$$

$$\text{Probability it is sunny on day 4: } S^4 = \frac{2}{3} [1 - (-0.2)^4] = \frac{2}{3} [1 - 0.0016] \approx \frac{2}{3}$$

(information about initial condition is lost very quickly)

5. (12pts) Consider diffusion in a thin cylindrical tube of constant cross-section radius R , with molecule loss through the side surface of the tube satisfying the following property: the loss per unit time per surface *side area* of the tube is proportional to concentration at a particular location, $u(x, t)$, with the constant of proportionality denoted γ . Derive the diffusion equation for the concentration $u(x, t)$ in this case. Start by re-deriving the conservation law.

$$\begin{aligned} \frac{dN(t)}{dt} &= A\Delta x \frac{\partial u(x^*, t)}{\partial t} \\ &= \left[\begin{array}{c} \text{Inflow rate} \\ \text{from the left} \end{array} \right] - \left[\begin{array}{c} \text{Outflow rate} \\ \text{from the right} \end{array} \right] - \left[\begin{array}{c} \text{loss} \\ \text{from the sides} \end{array} \right] = AJ(x, t) - AJ(x + \Delta x, t) - \gamma u(x^*, t) A_{SIDE} \end{aligned}$$

where $A_{SIDE} = 2\pi R\Delta x$ is the side area of a cylindrical "slice" (different compared to cross-section area $A = \pi R^2$)

$$\text{Now, divide by } \Delta x \text{ times } A: \quad \frac{\partial u(x^*, t)}{\partial t} = -\frac{J(x + \Delta x, t) - J(x, t)}{\Delta x} - \gamma \frac{A_{SIDE}}{A\Delta x} u(x, t)$$

$$\text{Limit } \Delta x \rightarrow 0 \Rightarrow \frac{\partial u(x, t)}{\partial t} = -\frac{\partial J}{\partial x} - \gamma u \frac{A_{SIDE}}{A\Delta x} \Rightarrow \frac{\partial u}{\partial t} = -\frac{\partial J}{\partial x} - \gamma u \frac{2\pi R\Delta x}{\pi R^2\Delta x} \equiv -\frac{\partial J}{\partial x} - \frac{2\gamma}{R} u$$

$$\text{Now plug into this equation the Fick's law of diffusion: } J = -D \frac{\partial u}{\partial x} \Rightarrow \boxed{\frac{\partial u}{\partial t} = D \frac{\partial^2 u}{\partial x^2} - \alpha u} \text{ where } \alpha = \frac{2\gamma}{R}$$

6. (17pts) Consider the following ODE in \mathbb{R}^2 : $\frac{d\mathbf{r}}{dt} = \mathbf{u}(\mathbf{r}) = \begin{pmatrix} -y \\ x - y^3 \end{pmatrix}$ (i.e. $\frac{dx}{dt} = -y$; $\frac{dy}{dt} = x - y^3$)

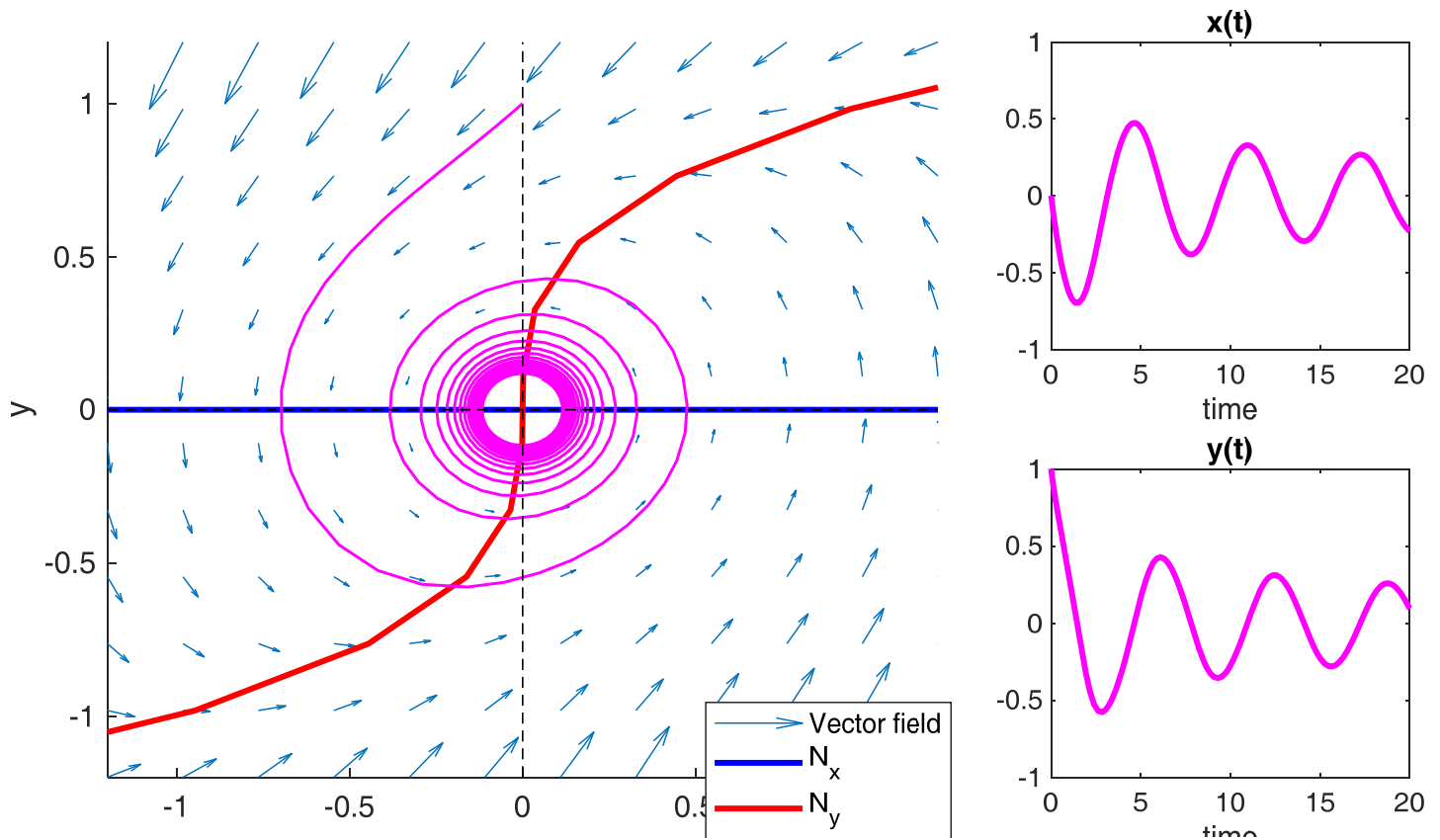
- Make a rough plot of the flow field in the (x, y) phase plane.
- Perform linear stability analysis of the equilibrium. Is linear stability analysis sufficient? Categorize the stability of the equilibrium.
- For the initial condition at $(0, 1)$, plot the trajectory in the (x, y) phase-plane, and plot $x(t)$ and $y(t)$ vs t . Be as accurate as possible.

$$J(\mathbf{r}_{eq}) = \begin{pmatrix} 0 & -1 \\ 1 & -y^2 \end{pmatrix} \Big|_{(0,0)} = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix} \Rightarrow \begin{cases} T=0 \\ D=1 \end{cases} \Rightarrow \lambda = \pm i$$

\Rightarrow Non-hyperbolic equilibrium \Rightarrow linear analysis is **insufficient**

$$\text{Examine distance from the origin: } \frac{d}{dt} r^2 = 2 \left(x \frac{dx}{dt} + y \frac{dy}{dt} \right) = 2(-xy + xy - y^4) = -2y^4 \leq 0$$

The origin is **asymptotically** stable (non-hyperbolic stable spiral)



7. (17pts) Consider the traffic flow equation, with physical traffic velocity depending linearly on traffic density:

$$\begin{cases} \frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}[(1-\rho)\rho] = 0 & (t > 0, x \in \mathbb{R}) \\ \rho(x, 0) \equiv \rho_0(x) = \begin{cases} -x, & x < 0 \\ 0, & x \geq 0 \end{cases} \end{cases}$$

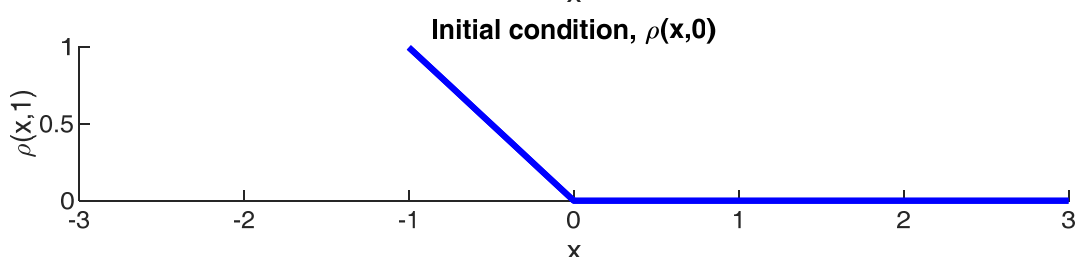
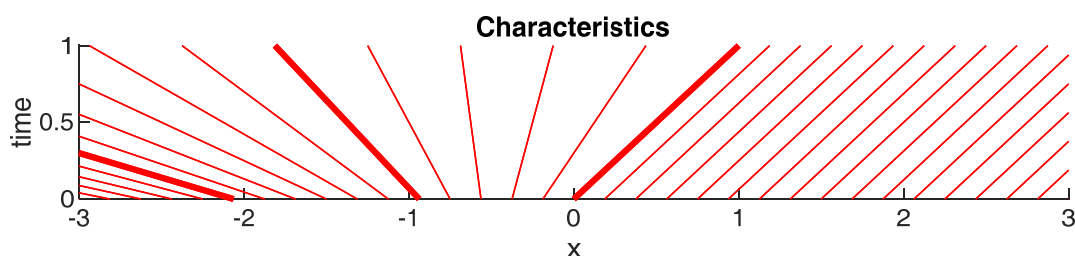
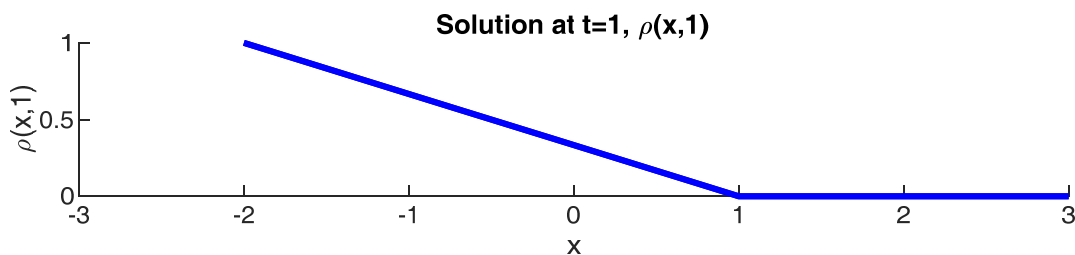
Don't forget to use the chain rule to convert this equation to standard advection form!

- Start by plotting the initial condition. Be careful with the minus signs.
- Plot the characteristics corresponding to $x_0 = -2$, $x_0 = -1$, and $x_0 = 0$. Is there a shock wave / break-up?
- Make a rough plot of traffic density $\rho(x, t)$ at $t=1$.
- Write down the explicit solution to this problem. It may help to separate the (x, t) domain into two regions.

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x}(\rho - \rho^2) = \frac{\partial \rho}{\partial t} + \underbrace{(1-2\rho)}_{\frac{dx}{dt}|_{cc}} \frac{\partial \rho}{\partial x} = 0 \Rightarrow \frac{dx}{dt}|_{cc} = 1-2\rho$$

$$x(t; x_0)|_{cc} = x_0 + (1-2\rho_0) t = \begin{cases} x_0 + (1+2x_0)t = (1+2t)x_0 + t, & x_0 < 0 \\ x_0 + t, & x_0 \geq 0 \end{cases} \Rightarrow \begin{cases} x_0 = -2: & x(t) = -2-3t \\ x_0 = -1: & x(t) = -1-t \\ x_0 = 0: & x(t) = t \end{cases}$$

$$\text{Invert: } x_0 = \begin{cases} \frac{x-t}{1+2t}, & x < t \\ x-t, & x \geq t \end{cases} \Rightarrow \rho(x, t) = \rho_0(x_0) = \begin{cases} \frac{t-x}{1+2t}, & x < t \\ 0, & x \geq t \end{cases} \Rightarrow \rho(x, 1) = \begin{cases} \frac{1-x}{3}, & x < 1 \\ 0, & x \geq 1 \end{cases}$$



8. (18pts) Convert to index notation, then use index notation to expand or simplify, and finally convert the result back to vector notation. Here $\mathbf{u}(\mathbf{r})$ is a smooth vector field, \mathbf{r} is the position vector, and $r = |\mathbf{r}|$:

a) $\nabla \times [\mathbf{r} \times \mathbf{u}(\mathbf{r})]$ Solved in last year's exam:

$$\begin{aligned} \nabla \times (\mathbf{r} \times \mathbf{u}) &= \varepsilon_{ijk} \partial_j (\mathbf{r} \times \mathbf{u})_k = \varepsilon_{ijk} \partial_j (\varepsilon_{knm} x_n u_m) = (\delta_{in} \delta_{jm} - \delta_{im} \delta_{jn}) \partial_j (x_n u_m) \\ &= \partial_m (x_i u_m) - \partial_n (x_n u_i) = x_i \partial_m u_m + u_m \underbrace{\partial_m x_i}_{\delta_{mi}} - x_n \partial_n u_i - u_i \underbrace{\partial_n x_n}_{\delta_{nn}=3} \\ &= x_i \partial_m u_m + u_i - x_n \partial_n u_i - 3u_i = \boxed{\mathbf{r} \cdot \nabla \cdot \mathbf{u} - (\mathbf{r} \cdot \nabla) \mathbf{u} - 2\mathbf{u}} \end{aligned}$$

b) $\nabla^2 \left(\frac{1}{r^p} \right)$, where $p = \text{const}$. For which p does this equal zero?

This was part of the calculation in problem 2 of homework 11, where you showed that the Laplacian of $1/r$ is zero. This is also why the potential of a point charge equals const / r (in other words, the Green's function of the Laplacian in free space = const / r). This gives the significance of this exercise

Use the chain rule: $\partial_m f(r) = f'(r) \frac{x_m}{r}$

$$\partial_m \partial_m \left(\frac{1}{r^p} \right) = \partial_m \left(-\frac{p}{r^{p+1}} \frac{x_m}{r} \right) = -p \left[x_m \partial_m \left(\frac{1}{r^{p+2}} \right) + \frac{1}{r^{p+2}} \partial_m x_m \right] = -p \left[-(p+2) \frac{\overbrace{x_m x_m}^{r^2}}{r^{p+4}} + \frac{\overbrace{\delta_{mm}}^N}{r^{p+2}} \right]$$

Now, assume $\mathbf{r} \in \mathbb{R}^N$: $= p \frac{p+2-N}{r^{p+2}} = 0$ if $p=0$ or $p=N-2$

If $N=3$: $p=0$ or $p=1$ as expected!