



Low temperature homoepitaxy of GaN by LP-MOVPE using Dimethylhydrazine and nitrogen

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Abstract

Thin films of GaN with the V/III \approx 10 ratio were grown by low-pressure metal organic vapour phase epitaxy (LP-MOVPE) using N₂ and Dimethylhydrazine (DMHy) as a carrier gas and nitrogen precursor, respectively. For the growth temperatures in the range from 550 to 690 °C the GaN layers exhibited good surface morphology. In the temperature range from approximately 550 to 610 °C, the growth rate increases with increasing temperature, characteristic of the process limited by surface kinetics with the activation energy of approximately 36 kcal/mol. For the temperatures between 620 and 690 °C, the growth rate was nearly independent of temperature, which is indicative of a mass transport limited growth. The activation energy was about 4.6 kcal/mol. Micro Raman spectroscopy revealed a significant relaxation of the selection rules for the scattering by the optical phonons in the films grown at lower temperatures. Variation of the intensity ratio for E₂^H and E₁ phonon modes has been attributed to the changes in the structural quality of the films grown at different temperatures.

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1. Introduction

Wide Band Gap semiconductors, such as GaN and related alloys (GaAlN, InAlN, InGaN), exhibit many attractive electronic and optical properties which make them ideal candidates for light emitting diodes (LEDs) and laser diodes operating throughout ultra violet to green visible spectral range as well as for high power and high frequency applications. One of the limiting aspects of GaN-based materials is their relatively poor crystal quality due to the necessity of heteroepitaxy. In some cases GaN buffer layer grown at low temperature on sapphire, silicon carbide or zinc oxide substrates is a necessary pathway process in order to improve the structural quality of these heterostructures. Ammonia (NH₃) is a conventional atomic nitrogen precursor for high temperature metal organic vapour phase epitaxy (MOVPE) growth of GaN-based materials. Unfortunately, this gas is not convenient for low temperature growth because of its low dissociation efficiency even at high temperature (only 15% decomposes at 950 °C) [1], which is attributed to the strong boundary energy N–H (104 kcal/mol). As a result, alternative precursors have been studied as hydrazine (N₂H₄) [2], Dimethylhydrazine ((CH₃)₂NNH₂) [3–7], and tertibutylhydrazine (tbu-NH₂) [8]. In addition to low temperature decomposition the nitrogen precursor has to fulfill more practical requirements such as low toxicity, high purity, liquid at room temperature and high vapour pressure.

In this work, we have chosen the DMHy as a nitrogen source, which has been used successfully in the deposition of GaN on GaAs [7], sapphire [3,5], Si(111) [4], GaN [6]. Our choice has the following advantages: (a) a relatively high dissociation efficiency at low temperature (50% decomposes at 420 °C) [9], (b) it is more stable than hydrazine, commercially available with high purity, (c) DMHy is liquid at room temperature and, finally, (d) it has high vapour pressure.

In order to explore GaN growth in wide range of low temperatures 550–690 °C, in this paper we combined DMHy with molecular nitrogen (N₂), an alternative carrier gas for hydrogen with a low thermal conductivity. In addition, this combination results in a uniform effective V/III ratio over the growth surface.

2. Experimental

The growths were carried out by MOVPE in a *T*-shape reactor [10] on *n*-type doped GaN wafers, obtained by deposition of 4 μm of GaN layers on sapphire substrates. The growth temperature ranged from 550 to 690 °C and was measured by a thermocouple placed in the graphite susceptor at few millimetres below the substrate. The growth pressure was 450 Torr, a process gas was nitrogen N₂, and the V/III ratio was equal to 11. TriMethylGallium (TMG) and Dimethylhydrazine (DMHy) were used as sources of gallium and nitrogen respectively. For all the growths the typical thickness layer was 350 nm. The growth rate was estimated *in situ* by reflectometry using a semiconductor laser operating at a wavelength of 670 nm. The laser beam was focused on the sample surface through a quartz reactor window and the spot size on the sample was 1 mm.

The surface morphology of the layers was examined by a combination of optical microscopy in the Nomarsky mode and atomic force microscopy.

The Raman spectra were measured with a Dilor confocal micro-Raman spectrometer using the excitation laser wavelength of 514.5 nm and the optical power at about 10 mW on the sample. The laser beam was focused through a ×100 numerical aperture objective lens. The confocal

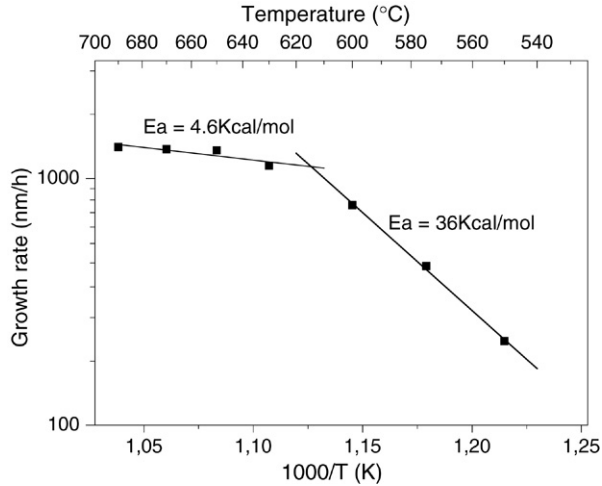


Fig. 1. Growth rate versus $1/T$ with the Arrhenius line fit.

aperture was fixed to analyze approximately $1 \mu\text{m}^3$ of the sampling volume. Polarization of the Raman spectra was analyzed in backscattering geometries.

3. Results and discussion

The growth rate was investigated as a function of the growth temperature (Fig. 1). In the temperature range from approximately 550 to 610 °C the growth rate increases with the temperature, characteristic of the process limited by surface kinetics. In other words, variation of the growth rate was related to the pyrolysis process of TMGa. The data points fit well using Arrhenius plot with the slope corresponding to the activation energy of approximately 36 kcal/mol. This value is slightly higher than the activation energy of 28.5 kcal/mol under H_2 carrier gas reported by Hsu et al. [6]. Between 620 and 690 °C, the growth rate has a much smaller slope corresponding to the activation energy of approximately 4.6 kcal/mol. This growth temperature regime is indicative of a mass transport limited growth. A similar trend in the growth rate vs. temperature was obtained in Ref. [13] with TMGa and NH_3 precursors.

Fig. 2 shows several time-dependencies of reflectivity measured for GaN layers grown at different temperatures. We have fixed arbitrarily the point zero in order to facilitate comparison between the reflectance curves. Their exact shape will be discussed below.

Visually, the surface of GaN layers was transparent with brownish colour. Under optical microscope with magnification of up to $\times 1000$, the substrate and all layers show the same surface morphology (Fig. 3). Below we present the materials quality studies using AFM and Raman spectroscopy.

Surface roughness has been characterized by the AFM (Fig. 4). At low temperature growth, a few islands appear on the surface of the layers. Above 600 °C, these islands are more numerous and more regular. The roughness and the z -average of the layers increase with the growth temperature (Table 1). The increase of RMS roughness for the layers is correlated to the decreases of reflectance intensity collected during the deposition (see Fig. 2). This phenomenon has been reported in the literature [14,15] and two explanations have been proposed (a) if the amplitude of reflectance oscillations is constant at the beginning of the growth and then decreases after few minutes, this situation indicates that GaN grows directly in 2D mode from the start. But the

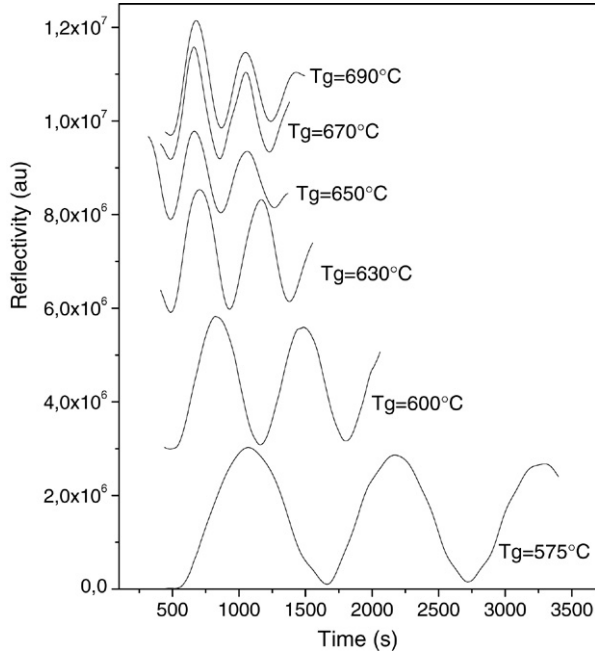


Fig. 2. Reflectance traces of growth versus the growth time.

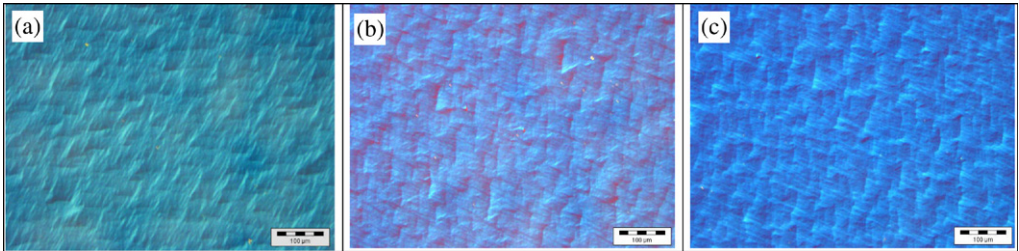


Fig. 3. Optical microscopy images of the GaN substrate (a), the layers grown at 575 (b) and 650 °C (c).

surface morphology is deteriorated by hillocks formation [14], (b) if this amplitude decreases right after the start of the growth, that would be a sign of 3D mode growth, and the average roughness increases with the island size. After the coalescence of the islands the reflectance intensity increases [15]. In our case, the amplitude of reflectance decreases just after the start of the deposition. Thus we can argue that the growths begin with formation of the 3D islands.

Raman characterizations: Gallium nitride on sapphire substrate crystallizes in hexagonal (wurtzite; space group C_{6v}^4) structure and the group theory predicts eight sets of phonon normal modes at the Γ point: $2A_1 + 2E_1 + 2B_1 + 2E_2$. The A_1 and E_1 modes are acoustic, while the other six modes, $A_1 + E_1 + 2B_1$ (silent) + $2E_2$ are optical [11]. Table 2 summarizes the Raman-active modes for gallium nitride in the backscattering $x(-, -)\bar{x}$ geometries [12]. The A_1 mode can be observed when the incident and scattered light have parallel polarization (e.g. $x(y, y)\bar{x}$ geometry), while the E_1 mode is observed only in crossed polarization geometry (e.g. $x(y, z)\bar{x}$). The E_2^H mode is allowed only in the $x(y, y)\bar{x}$ geometry.

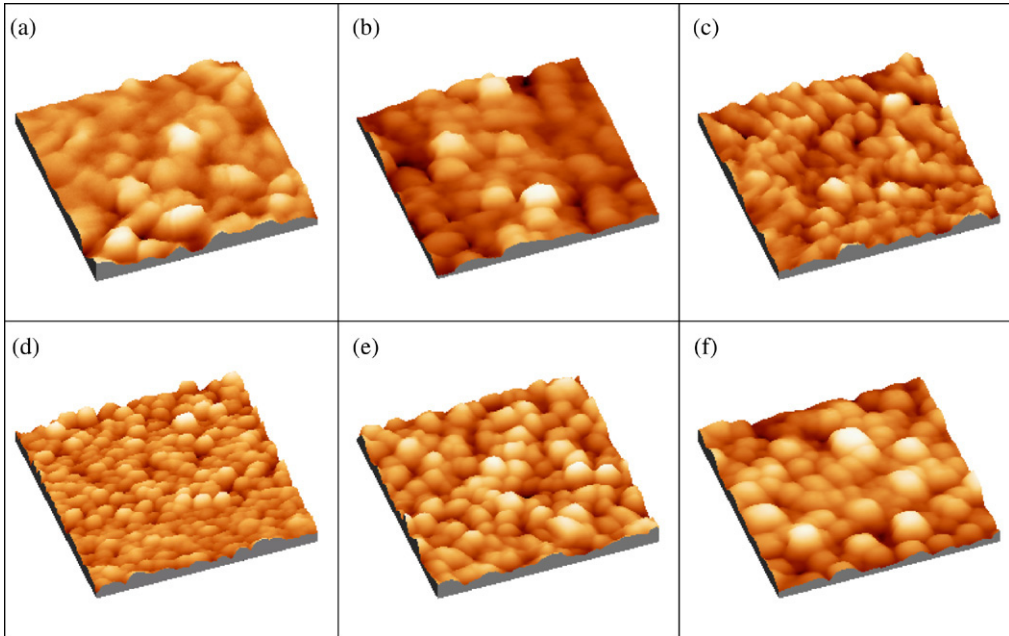


Fig. 4. AFM scans ($1 \mu\text{m} \times 1 \mu\text{m}$) taken on the layers obtained at 550 °C (a), 575 °C (b), 600 °C (c), 650 °C (d), 670 °C (e) and 690 °C (f).

Table 1
Evolution of RMS roughness and z average versus the growth temperature

Growth temperature (°C)	RMS roughness (nm)	z_{average} (nm)	z_{range} (nm)
550	1.47	5.75	11
575	2.05	6.5	15
600	2.25	7.5	17
650	2.9	11.4	24
670	4.3	14.5	28
690	4.4	13.6	28

Table 2
Raman-active modes in the $x(-, -)\bar{x}$ geometries [12]

Geometry	Raman allowed modes	Raman shift (cm^{-1})
$x(y, y)\bar{x}$	$A_1(\text{TO}), E_2^H$	533, 570
$x(y, z)\bar{x}$	$E_1(\text{TO})$	562
$x(z, z)\bar{x}$	$A_1(\text{TO})$	533

Fig. 5 shows evolution of the Raman spectra measured in $x(y, z)\bar{x}$ geometry for several samples grown at different temperatures. Raman selection rules allow only $E_1(\text{TO})$ modes in this configuration. However, we observed the E_2^H peak with a relative intensity (with respect to the $E_1(\text{TO})$ allowed peak) increasing with decreasing growth temperature. Fig. 6 shows the $E_2^H/E_1(\text{TO})$ intensity ratio as a function of the sample growth temperature. We note also the

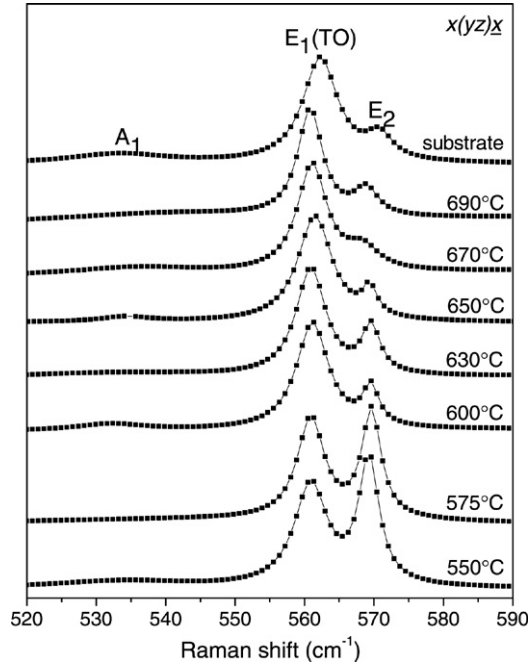


Fig. 5. The growth temperature dependence of Raman spectra in the $x(y, z)\bar{x}$ geometry. All spectra are normalized to the $E_1(\text{TO})$ peak intensity.

weak appearance of the $A_1(\text{TO})$ mode, which is forbidden in this configuration, with much lower relative intensity than the E_2^H one. The $E_2^H/E_1(\text{TO})$ peak intensity ratio measured for high-temperature-grown gallium nitride substrate is shown as a reference. The observed relaxation of the Raman selection rules can have two origins: (a) large solid angle of the objective lens of our Raman microscope [12] and (b) disorder in epitaxial layers. For our reference sample – high temperature grown GaN on sapphire – the first reason is dominant for observation of the forbidden modes in $x(y, z)\bar{x}$ geometry with the $E_2^H/E_1(\text{TO})$ peak intensity ratio of about 0.2. For the low temperature grown GaN samples relaxation of the Raman selection rules is significantly enhanced (Fig. 6) bringing the $E_2^H/E_1(\text{TO})$ peak intensity ratio above 1. It can only be explained by the crystalline properties of the grown material, namely by the growth temperature-related disorder in the epitaxial layers. Low-temperature samples exhibit the maximum degree of disorder, while for the samples grown at $T \geq 650^\circ\text{C}$ $E_2^H/E_1(\text{TO})$ peak intensity ratio became comparable to the reference GaN substrate, thus showing an increase in the layer quality. We can correlate this evolution to the morphology showed in the AFM images (Fig. 4): appearance of islands, probably crystallites at low growth temperature, which coalesce at higher temperature.

4. Conclusion

In summary, GaN materials have been grown at low temperature by employing DMHY and N_2 as precursor of nitrogen and carrier gas respectively. Surface roughness at different growth temperatures has been investigated by combining in situ reflectometry and AFM analysis. Micro Raman measurement underlined the enhancement of the structural quality and the degree of disorder as a function of growth temperature.

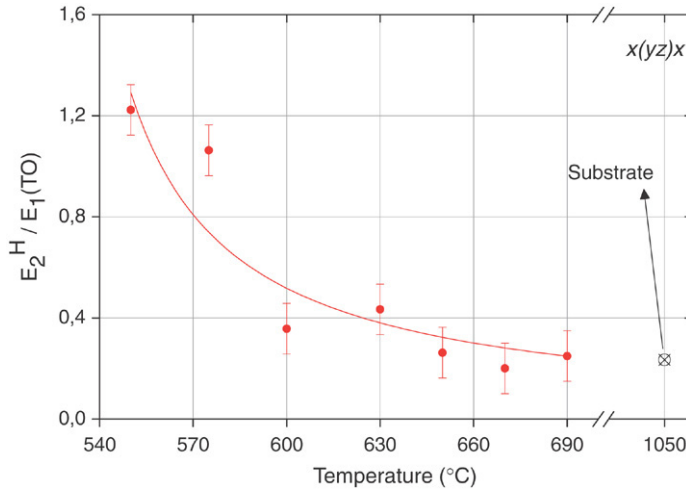


Fig. 6. Evolution of the Raman $E_2^H / E_1(TO)$ peak intensity ratio on the growth temperature. For comparison the ratio for the high temperature growth gallium nitride is shown (for this reference sample, used as substrate, the growth temperature is about 1050 °C).

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