

# Exploring the free energy landscape of a model $\beta$ -hairpin peptide and its isoform

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## ABSTRACT

Secondary structural transitions from  $\alpha$ -helix to  $\beta$ -sheet conformations are observed in several misfolding diseases including Alzheimer's and Parkinson's. Determining factors contributing favorably to the formation of each of these secondary structures is therefore essential to better understand these disease states.  $\beta$ -hairpin peptides form basic components of anti-parallel  $\beta$ -sheets and are suitable model systems for characterizing the fundamental forces stabilizing  $\beta$ -sheets in fibrillar structures. In this study, we explore the free energy landscape of the model  $\beta$ -hairpin peptide GB1 and its E2 isoform that preferentially adopts  $\alpha$ -helical conformations at ambient conditions. Umbrella sampling simulations using all-atom models and explicit solvent are performed over a large range of end-to-end distances. Our results show the strong preference of GB1 and the E2 isoform for  $\beta$ -hairpin and  $\alpha$ -helical conformations, respectively, consistent with previous studies. We show that the unfolded states of GB1 are largely populated by misfolded  $\beta$ -hairpin structures which differ from each other in the position of the  $\beta$ -turn. We discuss the energetic factors contributing favorably to the formation of  $\alpha$ -helix and  $\beta$ -hairpin conformations in these peptides and highlight the energetic role of hydrogen bonds and non-bonded interactions.

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**Key words:** potential of mean force; GB1; misfolded  $\beta$ -hairpin; hydrogen bonding; reptation.

## INTRODUCTION

Proteins fold into specific conformations, dictated by their amino acid sequence, resulting in unique functions. A variety of human diseases such as Alzheimer's, Parkinson's, and type II diabetes are associated with the incorrect folding of amyloid proteins.<sup>1,2</sup> These misfolded protein structures aggregate into fibrils exhibiting a high content of  $\beta$ -strands arranged into parallel<sup>3–9</sup> or antiparallel<sup>10–15</sup> orientations. Thus, protein misfolding can involve the structural reorganization from coil and/or  $\alpha$ -helical conformations to  $\beta$ -strands under denaturing conditions<sup>16,17</sup> and disease states.<sup>8,18,19</sup> A variety of factors including point mutations,<sup>20–22</sup> changes in pH,<sup>23</sup> and partially denaturing conditions<sup>24</sup> have been shown to influence the rate of  $\alpha$ -helix to  $\beta$ -sheet transition.

Experiments suggest that association of misfolded proteins into fibrils occurs for all amino acid sequences under appropriate denaturing conditions, implying that fibril structures emerge from properties that are common to all proteins, that is, the backbone.<sup>1</sup> In this scenario, the role of side chain interactions is to modulate the propensity for fibril formation. Accordingly, experiments have shown that increasing the hydrophobicity of nucleating

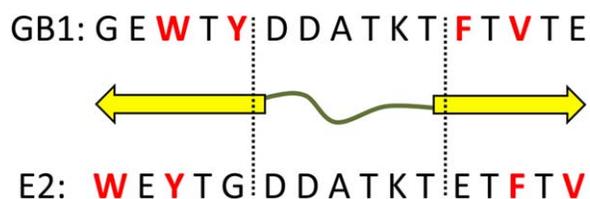
regions<sup>25,26</sup> and the presence of aromatic residues<sup>27</sup> increases the propensity of fibril formation. Moreover, hydrophobic interactions favored the formation of  $\beta$ -sheet structures<sup>28–30</sup> and they penalized  $\alpha$ -helical conformations in model peptides.<sup>29,30</sup> This could rationalize why some amyloid peptides form preferentially  $\beta$  or coil structures in aqueous environments and  $\alpha$ -helical conformations in membrane mimicking environments.<sup>31–33</sup> In contrast to hydrophobic interactions, mutations increasing the net charge of peptides inhibit fibril formation.<sup>34</sup> Computer simulations have also explored pathways for the inter-conversion of  $\alpha$ -helix to  $\beta$ -sheet. This inter-conversion occurs either through a random coil intermediate<sup>35</sup> or through a zipping mechanism.<sup>36</sup>

$\beta$ -Hairpin peptides are suitable model systems for characterizing the fundamental forces stabilizing

Additional Supporting Information may be found in the online version of this article.

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**Figure 1**

Primary sequence of GB1 and the E2 isoform and a schematic representation of the secondary structure corresponding to GB1. Hydrophobic residues are indicated in red. [Color figure can be viewed in the online issue, which is available at [wileyonlinelibrary.com](http://wileyonlinelibrary.com).]

$\beta$ -sheets. In particular, GB1 peptide (GEWTYDDATKTFVTE), the C-terminal  $\beta$ -hairpin region (residues 41–56) of protein G, is amenable to computational studies<sup>37–48</sup> due to its small size, stability of its  $\beta$ -hairpin conformation, and fast folding kinetics ( $\sim 6 \mu\text{s}$ ).<sup>26,38</sup> Garcia and Sanbonmatsu<sup>37</sup> first showed using replica exchange simulations in explicit water that GB1 has a small propensity for  $\alpha$ -helical conformations, in addition to the predominant  $\beta$ -hairpin conformation. The  $\alpha$ -helical conformations were also observed in subsequent studies of the peptide in implicit solvent.<sup>38,49</sup> Levy and coworkers<sup>49</sup> calculated the free energy difference between the  $\alpha$ -helical and  $\beta$ -hairpin states to be  $1.573 \pm 0.018$  kcal/mol. They showed that the transition from  $\alpha$ -helix to  $\beta$ -hairpin conformations begins with the association of the hydrophobic core even before the formation of native hydrogen bonds. Hence, changing the position of non-polar residues within the amino acid sequence is expected to affect the formation of the hydrophobic core and the stability of  $\alpha$  and  $\beta$  conformers.<sup>50–52</sup> In an interesting study, Ma and Nussinov<sup>40</sup> designed the E2 isoform of GB1 (WEYTGDDATKTETFTV) by moving the hydrophobic residues to the termini (see Fig. 1). Sequence-based structure prediction suggests that both peptides adopt  $\beta$ -hairpin conformation (see Supporting Information Fig. S1); however, Nussinov *et al.* showed that the energy landscape of the E2 isoform strongly favors helical conformations. This is remarkable as hydrophobic contacts in both peptides involve same residue pairs W–V and Y–F, and they have identical turn sequence (Fig. 1). It is an indication that  $\beta$ -sheet structures in GB1 are not formed solely by hydrophobic collapse or turn propensity but perhaps through a combination of both factors.<sup>53</sup>

In this study, we use all-atom simulations in explicit water to explore the free energy landscapes of GB1 and the E2 isoform. To this effect, we use umbrella sampling<sup>54</sup> with harmonic restraints applied to the ends of the chains to sample folded and unfolded conformations of the peptides. We find that the global minimum of the free energy for GB1 can be mapped to  $\beta$ -hairpin conformations whereas its unfolded state is populated with

misfolded hairpins differing from each other in the location of the  $\beta$ -turn. In contrast, the energy landscape of the E2 isoform is characterized by a large presence of helical conformations. We show that an increase in peptide–peptide and water–water hydrogen bonds is observed during the formation of these structures. However, this increase is compensated by a decrease in peptide–water hydrogen bonds resulting in an insignificant change in the net number of hydrogen bonds within the explored range of end-to-end distances. These results suggest that hydrogen bonds are not main interactions driving the formation of  $\alpha$ -helical and  $\beta$ -hairpin secondary structures, supported by the unfavorable electrostatic energy observed for the formation of these structures.

## SIMULATION METHODOLOGY

Initial structures were prepared from amino acid sequences of GB1 (GEWTYDDATKTFVTE) and the E2 isoform<sup>40</sup> (WEYTGDDATKTETFTV). Simulated annealing was performed on the extended conformations to create random structures of the peptides. Structures with the least number of hydrogen bonds were chosen from the simulated annealing step to avoid any bias of the starting conformation on the final structures.

We performed molecular dynamics simulations of the uncapped peptides at 310 K using GROMACS (Groningen MAchine Chemical Simulations, version 4.5.4)<sup>55</sup> simulation package and the AMBER ff99sb-ILDN<sup>56–58</sup> force field. Periodic boundary conditions were applied in all directions. To explore the  $\alpha$ -helical and  $\beta$ -hairpin conformations of GB1 and the E2 isoform, we used the pull code in GROMACS<sup>55</sup> to simulate the peptides over a range of end-to-end distances between 0.3 and 2.4 nm at 0.1 nm intervals. An additional simulation was performed at 0.45 nm, which corresponds to the end-to-end distance at which a perfect  $\beta$ -hairpin is observed. For each simulation, the end-to-end distance was restrained to a specific value within the above-mentioned range. A force constant of  $25 \text{ kcal/mol/\AA}^2$  was used for the pull simulations.  $\beta$ -Hairpin conformations are expected at lower end-to-end distances ( $\sim 0.5$  nm) while  $\alpha$ -helical conformations are expected at larger end-to-end distances ( $\sim 2.2$  nm).

Peptides were immersed in a  $4.5 \text{ nm} \times 4.5 \text{ nm} \times 4.5 \text{ nm}$  solvent box containing 2892 and 3623 TIP3P<sup>59</sup> water molecules for the GB1 and E2 isoform, respectively. Three  $\text{Na}^+$  ions were added to neutralize the system. Protein and non-protein atoms were coupled to their respective temperature baths and maintained at 310 K. Temperature and pressure were maintained CONSTANT using the v-rescale thermostat<sup>60</sup> and Parrinello-Rahman barostat,<sup>61</sup> respectively. Simulations were performed using a 2 fs timestep. Electrostatic interactions were calculated using the Particle-Mesh Ewald approach with a

grid spacing of 0.12 nm. Electrostatic cut-off of 1.3 nm and van der Waals cut-off of 1.2 nm were used. Two sets of independent simulations were performed over the specified end-to-end distance range corresponding to a total of 23 simulations for each set. Each of the end-to-end distance simulations was performed for 400 ns. We used the last 300 ns of each of the two independent sets of simulations over the entire end-to-end range of distances. Potential of mean force (PMF) was calculated as a function of the end-to-end distance using the weighted histogram analysis method (WHAM)<sup>62</sup> using `g_wham`<sup>63</sup> implemented in GROMACS.

Secondary structures were calculated using the DSSP algorithm<sup>64</sup> which determines hydrogen bonds using an electrostatic model. In Figure 2(A,B), a strand refers to either a parallel or an anti-parallel bridge while a helix refers to  $\alpha$ -helix,  $3_{10}$  helix, or  $\pi$ -helix. For the hydrogen bond analysis in Figure 2(E,F), we used a common used geometrical definition common geometry in which hydrogen bonds are formed when the distance between donor (D) and acceptor (A) is smaller than 0.35 nm and the angle A–D–H is smaller than 40°. We used `g_hbond` provided in the software package GROMACS for this calculation.

Representative conformations shown in Figure 2(A,B) were selected by clustering simulation ensembles. Clustering was performed using the `gromos` clustering method provided with the `g_cluster` utility of GROMACS.<sup>65</sup> Further details are provided in supporting Information. In Figure 2(E–H), rescaled values of hydrogen bonds and energetic terms (van der Waals, Coulomb, and total) were calculated as the difference of these quantities at each end-to-end distance and that at 2.4 nm. Plots of the actual means for these parameters (without rescaling) are presented in Supporting Information Figures S3 and S4.

## RESULTS AND DISCUSSION

### Free energy landscape of GB1 and the E2 isoform shows distinct preferences for $\beta$ -hairpin and $\alpha$ -helical conformations

The free energy landscape of GB1 and the E2 isoform is computed along the reaction coordinate corresponding to the peptide end-to-end distance in the range of 0.3–2.4 nm. Figure 2(A,B) shows the PMFs for the GB1 peptide and the E2 isoform at 310 K. Qualitative agreement is observed between the minima in the PMF and the secondary structural propensities of the peptides [Fig. 2(C,D)]. Global minimum of the free energy for GB1 is observed at end-to-end distances of 0.5 nm and it is characterized by structures with  $\beta$ -hairpin conformations.  $\beta$ -Hairpin propensity is prominent over an extended range of end-to-end distances from 0.4 to 1.6 nm [Fig. 2(C)]. Representative conformations corresponding to the different minima are shown in Figure

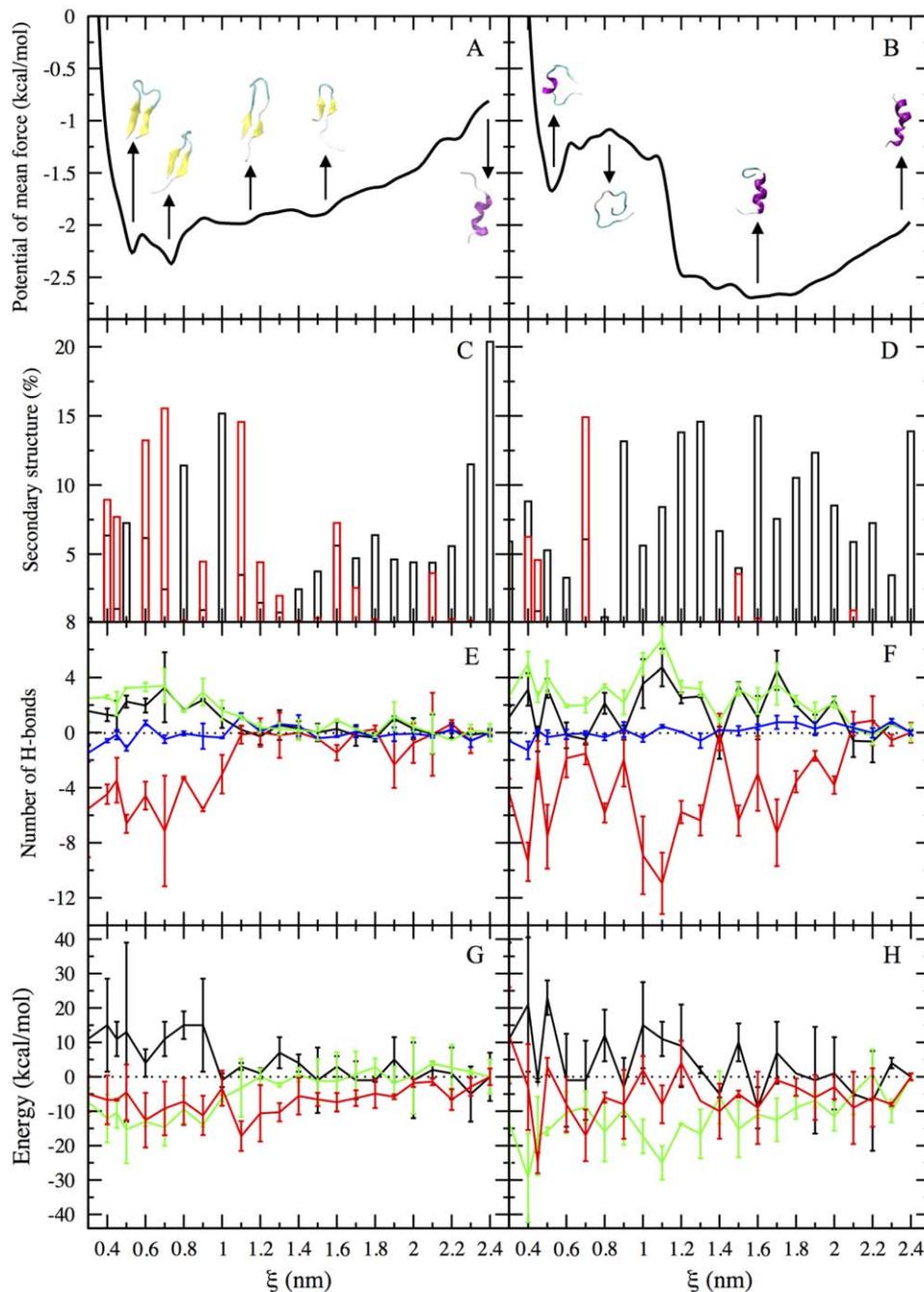
2(A). GB1 adopts a significant amount ( $\sim 20\%$ ) of completely folded helical conformations at larger end-to-end distances ( $\sim 2.4$  nm). However, these conformations are not energetically favored. The free energy difference between the  $\beta$ -hairpin and  $\alpha$ -helical conformations of GB1 is calculated to be  $\sim 1.56$  kcal/mol, consistent with that reported by Levy and coworkers.<sup>49</sup>

The free energy landscape of the E2 isoform is drastically different from that of the GB1 peptide. The global minimum of the free energy is observed at intermediate end-to-end-distances between 1.1 and 1.8 nm. These regions of free energy minima correspond to structures with partially folded  $\alpha$ -helical conformations [Fig. 2(D)]. Although this peptide has a small propensity for  $\beta$ -hairpin conformations, these structures are not energetically favorable. This strong preference of the E2 isoform for helical conformations is consistent with observations of Ma and Nussinov.<sup>40</sup> Further, completely folded helical structures (observed at  $\xi > 2.0$  nm) are less favorable compared to partially folded ( $1.2 \text{ nm} < \xi < 1.8 \text{ nm}$ ) helices. Partially folded helical conformations have also been observed in numerous helix stability and denaturation studies of helical peptides and proteins.<sup>35,66–68</sup>

### Misfolded $\beta$ -hairpin conformations of GB1

In the zipper-like mechanism for  $\beta$ -hairpin formation, native turn contacts form first followed by hydrogen bonds which zip the two strands together all the way to the N- and C-termini. During zipping, the position of the turn does not move as the end-to-end distance of the peptide decreases [Fig. 3(A)]. This symmetry is clearly not observed in our simulations of GB1 as can be seen from the characteristic configurations shown in Figure 2(A). We observe three misfolded  $\beta$ -hairpin conformations of GB1 differing from each other in the position of the turn. One of these conformations with non-native turn close to the N-terminus [similar to the represented structure at 1.6 nm in Fig. 2(A)] was observed previously using all-atom simulations in explicit water.<sup>69</sup> Furthermore, it was shown that the presence of this misfolded  $\beta$ -hairpin is consistent with the two state folding kinetics of GB1.<sup>69</sup> To the best of our knowledge, this is the first time that the other two misfolded  $\beta$ -hairpins shown in Figure 2(A) are reported.

In the misfolded  $\beta$ -hairpins, the turn is displaced from the N-terminus towards the middle of the chain with decreasing end-to-end distance, as shown in Figure 4. For large end-to-end distances ( $\xi = 1.6$  nm), turns can be found anywhere from the middle of the sequence to the C-terminus. As the end-to-end distance decreases, the turn becomes more localized closer to its native position (residues 7–10). The free energy landscape of GB1 does not show large energy barriers separating the different misfolded  $\beta$ -hairpins. This suggests that when one of the misfolded  $\beta$ -hairpins is formed, the system can reach

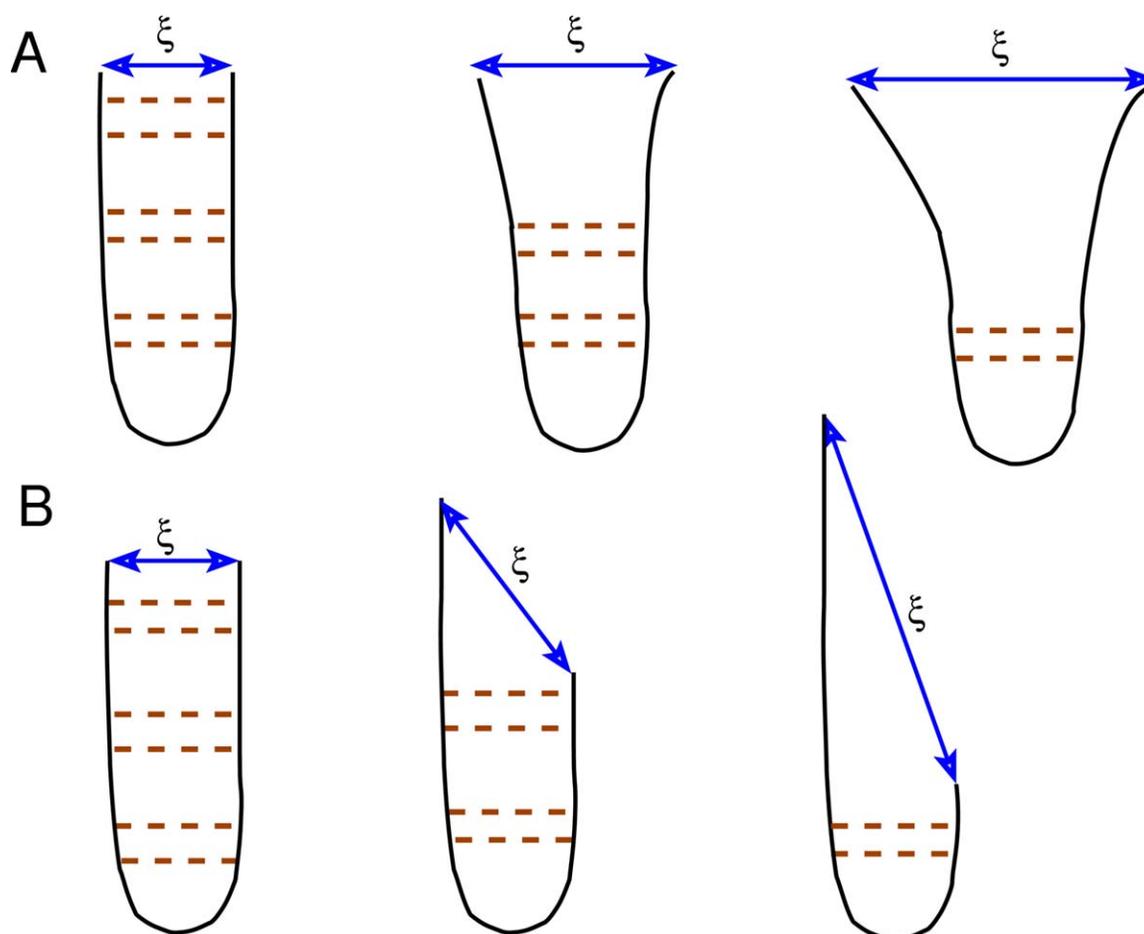


**Figure 2**

(A, B) PMF as a function of the end-to-end distance ( $\xi$ ) for the GB1 peptide (A) and the E2 isoform (B). Representative conformations corresponding to the different energy minima and the free energy barriers for the GB1 peptide and the E2 isoform are shown in a cartoon representation. (C, D) Percent secondary structure as a function of  $\xi$  for GB1 (C) and the E2 isoform (D). Helical content is shown in black while the  $\beta$ -strand content is displayed in red. (E, F) Rescaled values for the peptide-peptide (black), peptide-water (red), water-water (green), and total (blue) hydrogen bonds as a function of  $\xi$  for GB1 (E) and the E2 isoform (F). (G, H) Rescaled values for the Coulomb (black), van der Waals (green), and total (red) energies as a function of  $\xi$  for GB1 (G) and the E2 isoform (H). Error bars represent the standard error of the mean from the two simulations.

the other hairpins and the native structure through a reptation-like mechanism.<sup>47</sup> In this mechanism, hydrogen bonds break and reform during the transition from the different misfolded structures to the native-like hair-

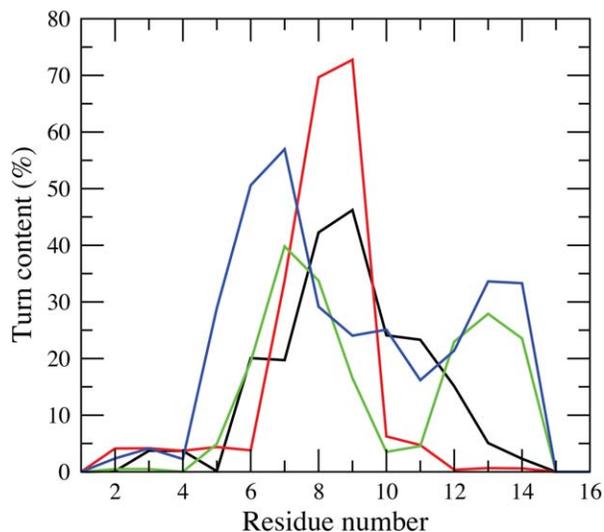
pin conformation [Fig. 3(B)]. Reptation mechanism has been reported previously<sup>47</sup> for GB1 although most computer simulations report  $\beta$ -hairpin formation initiated at the native turn.<sup>45,52,70,71</sup>

**Figure 3**

Schematic representation of the zipper-like (A) and roll-up (B) mechanisms for  $\beta$ -hairpin formation. [Color figure can be viewed in the online issue, which is available at [wileyonlinelibrary.com](http://wileyonlinelibrary.com).]

Misfolded  $\beta$ -hairpins could also emerge as the result of a roll-up mechanism proposed recently for Chignolin.<sup>72</sup> In this mechanism, one end of the peptide is kept rigid while the other rolls-up displacing the turn to the middle of the chain. The enhanced rigidity of one end with respect to the other is a key constituent for folding to proceed. Enemark *et al.*<sup>72</sup> showed that this rigidity is conferred by frequent interactions between aromatic residues at positions  $i$  and  $i + 2$  which restrict main-chain movement. This could be the rationale for the presence of two aromatic residues at a short distance from each other at the rigid N-termini of GB1. The presence of aromatic residues at positions  $i$  and  $i + 2$  is however not enough to guarantee a robust energy landscape favoring  $\beta$ -hairpins as demonstrated by the E2 isoform peptide. These aromatic residues have to be positioned close to the turn,<sup>40,50–52</sup> suggesting that the intrinsic propensity of residues 7–10 of GB1 to form turns is not enough to account for the native structure. It requires a reinforcement of hydrophobic interactions to be stabilized.

The misfolded  $\beta$ -hairpins observed in this study correspond to characteristic structures of equilibrium ensembles which have fixed end-to-end distances. Based on our simulations alone, it is, therefore, not possible to determine the sequential order of events that leads to correctly folded  $\beta$ -hairpin which is required to discriminate between reptation, zipper, and roll-up mechanisms. Nevertheless, our simulations predict that the ensemble of unfolded structures will be populated by misfolded  $\beta$ -hairpins. To assess if these structures are indeed observed in kinetic simulations, we performed a set of five unconstrained simulations of the GB1 peptide. Simulations were performed using starting conformations with end-to-end distances of 1.6 nm for a total of 300 ns. We identified misfolded conformations similar to those observed in the end-to-end restrained simulations (see Supporting Information Fig. S2), suggesting that misfolded  $\beta$ -hairpins observed in our simulations are not artifacts of the sampling methodology.



**Figure 4**

Average percent turn content as a function of the residue number for end-to-end distances of 0.4 (black), 0.7 (red), 1.2 (green), and 1.6 nm (blue) for the GB1 peptide. [Color figure can be viewed in the online issue, which is available at [wileyonlinelibrary.com](http://wileyonlinelibrary.com).]

#### Hydrogen bonds are not main interactions driving secondary structure formation

Figure 2(E,F) shows the peptide–peptide, peptide–water, water–water, and total hydrogen bonding calculated for GB1 (E) and the E2 isoform (F). For GB1, the number of peptide–peptide hydrogen bonds (black) increases at end-to-end distances corresponding to  $\beta$ -hairpin conformations. The number of water–water hydrogen bonds (green) shows a similar trend, suggesting that as the peptide forms intra-chain hydrogen bonds, the displaced water molecules are transferred to the bulk, resulting in a corresponding increase in the number of water–water hydrogen bonds. This is supported by the observation that the number of peptide–water hydrogen bonds (red) decreases as the peptide–peptide and water–water interactions increase. We note that the net number of hydrogen bonds (blue) does not change significantly for all values of  $\xi$ . This is in agreement with Schellman's<sup>73</sup> and Kauzmann's<sup>74</sup> hypotheses that hydrogen bonds in aqueous solution do not play an important role in protein folding.<sup>75</sup> Our results are also in agreement with previous simulations on *N*-methylacetamide<sup>76</sup> and homo-peptide dimers made of aliphatic side chains.<sup>77</sup> These simulations also reported negligible role of hydrogen bonding during secondary structure formation. Nevertheless, this is surprising because it contradicts most interpretations of the zipper mechanism for which hydrogen bonding is a key component in zipping the two strands together.

Hydrogen bonding trends for the E2 isoform are similar to that of the GB1 peptide. The number of peptide–

peptide hydrogen bonds (black) is higher for the preferred  $\alpha$ -helical conformations at intermediate end-to-end distances suggesting greater intra-peptide hydrogen bonding [Fig. 2(F)]. A corresponding increase in the water–water hydrogen bonding is observed at these end-to-end distances. The number of peptide–water interactions is lower for the regions where the peptide–peptide and water–water interactions are higher. The total number of hydrogen bonds (blue) does not show a significant change as a function of the end-to-end distance for the E2 isoform.

#### Coulomb and van der Waals energies contribute differently to the formation of secondary structures

Figure 2(G,H) shows the rescaled values (see Methods) of the average Coulomb (black), van der Waals (green), and total energies (red) as a function of the end-to-end distance for GB1 and the E2 isoform. Coulomb energies are large and unfavorable at small end-to-end distances ( $<1$  nm) and these values decrease with increasing end-to-end distance. In contrast, van der Waals interactions contribute favorably to the formation of  $\beta$ -hairpin structures. The total energy is dominated by the behavior of van der Waals interactions and it is negative along the entire reaction coordinate for GB1. This is different from results from implicit water simulations which suggest that electrostatic interactions are main driving forces for folding of GB1.<sup>78</sup>

Energetic contributions for the E2 isoform are subjected to larger fluctuations. The average total energy for the E2 isoform is favorable along entire length of the reaction coordinate, similar to that observed for the GB1 peptide. Coulomb interactions are unfavorable for most of the end-to-end distances explored while van der Waals interactions are favorable over the entire range of end-to-end distances.

## CONCLUSIONS

Transitions from  $\alpha$ -helix or coil states to  $\beta$ -sheets are essential for the formation of fibrillar aggregates, observed in various disease states.<sup>19,23,36,79,80</sup> Here, we used the GB1 peptide and the E2 isoform as model peptides to study large conformational changes associated with the formation of  $\beta$ -hairpin and  $\alpha$ -helical structures. Peptides were simulated over a range of end-to-end distances that explore both  $\beta$ -hairpin ( $\sim 0.4$ – $0.7$  nm) and  $\alpha$ -helical ( $\sim 2.2$ – $2.4$  nm) conformations. We examined free energy landscapes of GB1 and the E2 isoform in explicit solvent. In agreement with previous studies<sup>37,49</sup> we show that  $\beta$ -hairpin is the minimum free energy conformation for GB1. In addition to the presence of helical conformations close to the C-termini,<sup>49</sup> the unfolded state of GB1 is also populated by misfolded  $\beta$ -hairpins

which differ from each other in the location of the  $\beta$ -turn. The existence of one of the misfolded  $\beta$ -hairpins was reported previously.<sup>69</sup> These observations suggest that the unfolded state of GB1 may not be as disordered as commonly thought. In contrast, the E2 isoform adopts predominantly  $\alpha$ -helical conformations with a stable partially folded  $\alpha$ -helix being the preferred free energy minimum.<sup>40</sup>

In a recent computer simulation, side chain interactions were shown to be a main force driving  $\beta$ -sheet formation of model peptides while hydrogen bonds were shown to play a minimal role in this process.<sup>77</sup> A limitation of this work was that peptides were kept elongated during the simulations. Here, we show that even when the backbone of the peptide is not restrained, the total number of hydrogen bonds does not change during the formation of the  $\beta$ -hairpin while the total electrostatic energy of the system becomes unfavorable, that is, it increases. Hydrogen bond energies in all-atom models are a result of electrostatic interactions between hydrogen donor and acceptor groups. This suggests that hydrogen bonding might not be a major driving force for the formation of secondary structures in GB1 and the E2 isoform. A possible role of hydrogen bonding might be to restrict the conformational space of proteins, penalizing structures for which the total number of hydrogen bonds is not optimized. Configurations where hydrogen bond donors or acceptors of the backbone are unsatisfied would have a large electrostatic energy and would therefore occur rarely. We hypothesize that restriction of the conformational space to configurations with optimized hydrogen bonding could be a key factor accounting for fibril formation.

From an enthalpic point of view, the formation of  $\beta$ -sheets in the GB1 peptide is characterized by favorable van der Waals interactions and unfavorable Coulomb energies. This implies that as far as Coulomb interactions are concerned, extended conformations with large peptide exposure to water are more favorable than  $\beta$ -hairpin conformations. In contrast, van der Waals interactions favor the more compact  $\beta$ -hairpin conformation. Although average values of the different energetic terms computed for the E2 isomer are subjected to large fluctuations in our simulations, they exhibit a similar trend: compact conformations have favorable van der Waals and unfavorable Coulomb energies.

Entropic energies ( $-T\Delta S$ ) at different end-to-end distances for both systems studied here were not computed due to their prohibitive computational cost. However, we can speculate that they have two main contributions which are due to water molecules around non-polar residues, that is, shell water, and the backbone of the peptide. Shell water around non-polar residues have a high Gibbs free energy (low entropy) due to the reduced number of conformations they assume compared to bulk water. Thus, entropic energies due to shell water favor

compact conformation with small end-to-end distances as these conformations have a smaller number of shell water. In contrast, the backbone of the peptide assumes a larger number of conformations in the unfolded state which is, therefore, favored by backbone entropy.

In summary, the energetic picture that emerges from this study is that Van der Waals and hydrophobic interactions drive the formation of  $\beta$ -hairpins in the GB1 peptide while Coulomb interactions and the entropy of the backbone oppose it. A similar picture also emerges for the E2 isomer. However, our simulations for the E2 isomer are subjected to large error bars.

The objective of this study was to explore the free energy landscape associated with the transition from  $\beta$ -hairpin to  $\alpha$ -helical conformations. The end-to-end distance of the peptide was chosen as the reaction coordinate since (1) it distinguishes  $\alpha$ -helices from  $\beta$ -hairpin conformations, and (2) it is suitable for umbrella sampling simulations. Although the end-to-end distance is an unnatural constraint, structures obtained in this study, for example, misfolded  $\beta$ -hairpin of GB1, are also observed in unconstrained simulations of GB1, indicating that conformations obtained from the constrained simulations are not artifacts. Further, this study can be extended to explore energy landscapes of naturally occurring chameleon peptides<sup>81,82</sup> that have an equal propensity to adopt both  $\alpha$ -helix and  $\beta$ -strand conformations. This methodology can also be used effectively to characterize structural transitions from  $\alpha$ -helix to  $\beta$ -sheet conformations observed in amyloid- $\beta$  and other peptides involved in disease states.<sup>19,23,36,79,80</sup>

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