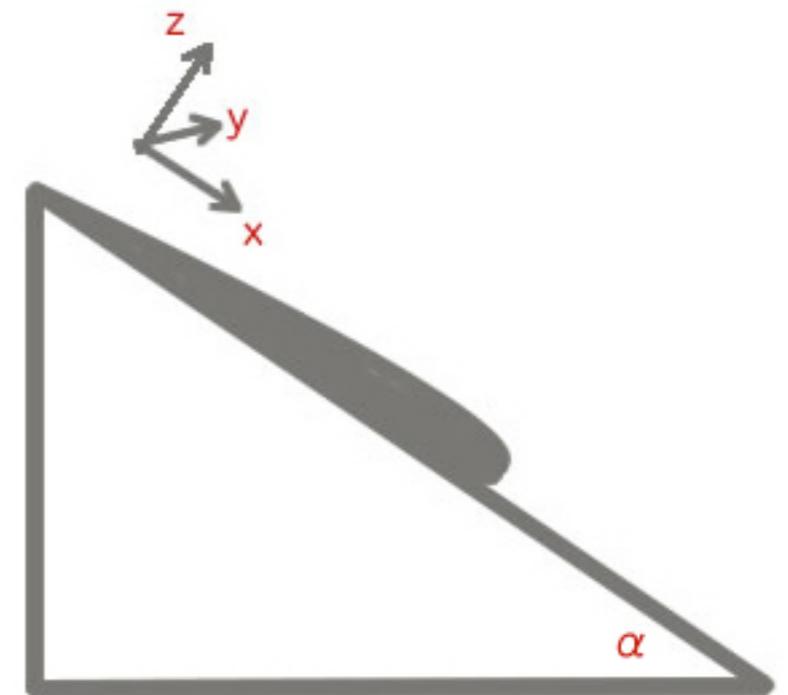


Nematic liquid crystals flowing down an incline

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Lubrication theory

- Thin film of a viscous fluid
- lubrication approximation
 - $\epsilon = H / L$ where $H \ll L$
 - velocity gradients in the x, y -directions negligible compared to velocity gradient in z -direction
 - $\epsilon^2 Re \ll 1$, can ignore terms due to inertia
- Lubrication theory simplifies evolution equation

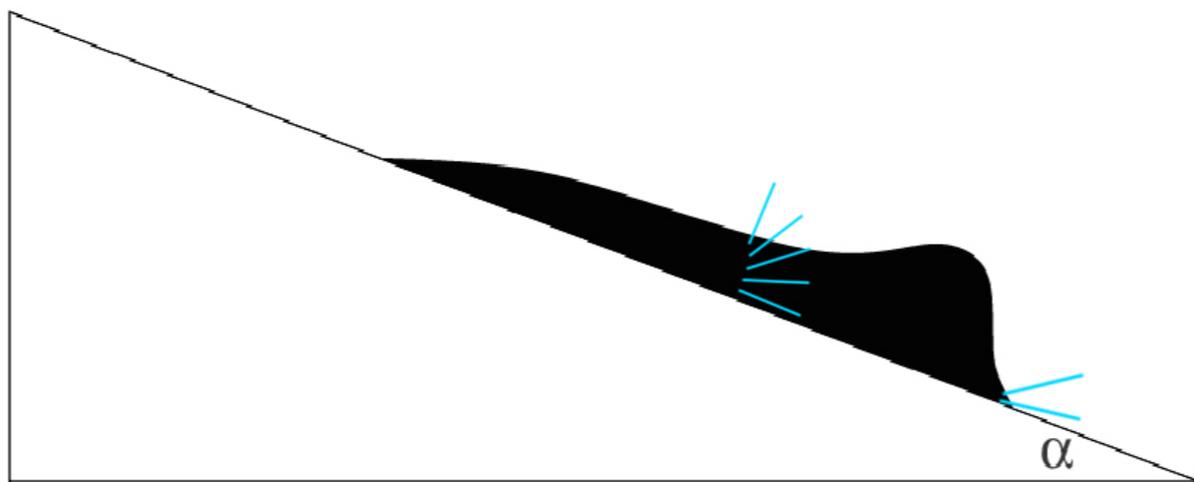


Evolution equation

$$h_t + \nabla \cdot [h^3 (C \nabla \nabla^2 h - B \nabla h) + N(m^2 - h m m') \nabla h] + U(h^3)_x = 0$$

$$m = \frac{h^{3/2}}{\beta^{3/2} + h^{3/2}}$$

- lubrication approximation adopted
- prewetting of substrate with precursory layer of thickness $b \ll h$
- weak-anchoring model



h : fluid thickness
 C : inverse capillary number
 B : Bond number,
 N : inverse Ericksen number

Linear stability analysis(LSA) of flat film

- Assume h is independent of y

$$h_t + \partial_x [h^3 (Ch_{xxx} - Bh_x) + N(m^2 - h m m') h_x] + U(h^3)_x = 0$$

- Perturb profile by a small amplitude

$$h(x,t) = h_o + \epsilon h_1(x,t) + O(\epsilon^2)$$

$$h_{1t} + Ch_o^3 h_{1xxxx} - (B - NM(h_o)) h_o^3 h_{1xx} + 3U h_o^2 h_{1x} = 0$$

$$M(h_o) = \frac{h_o^{3/2} - \beta^{3/2} / 2}{(h_o^{3/2} + \beta^{3/2})^3}$$

LSA of flat film cont.

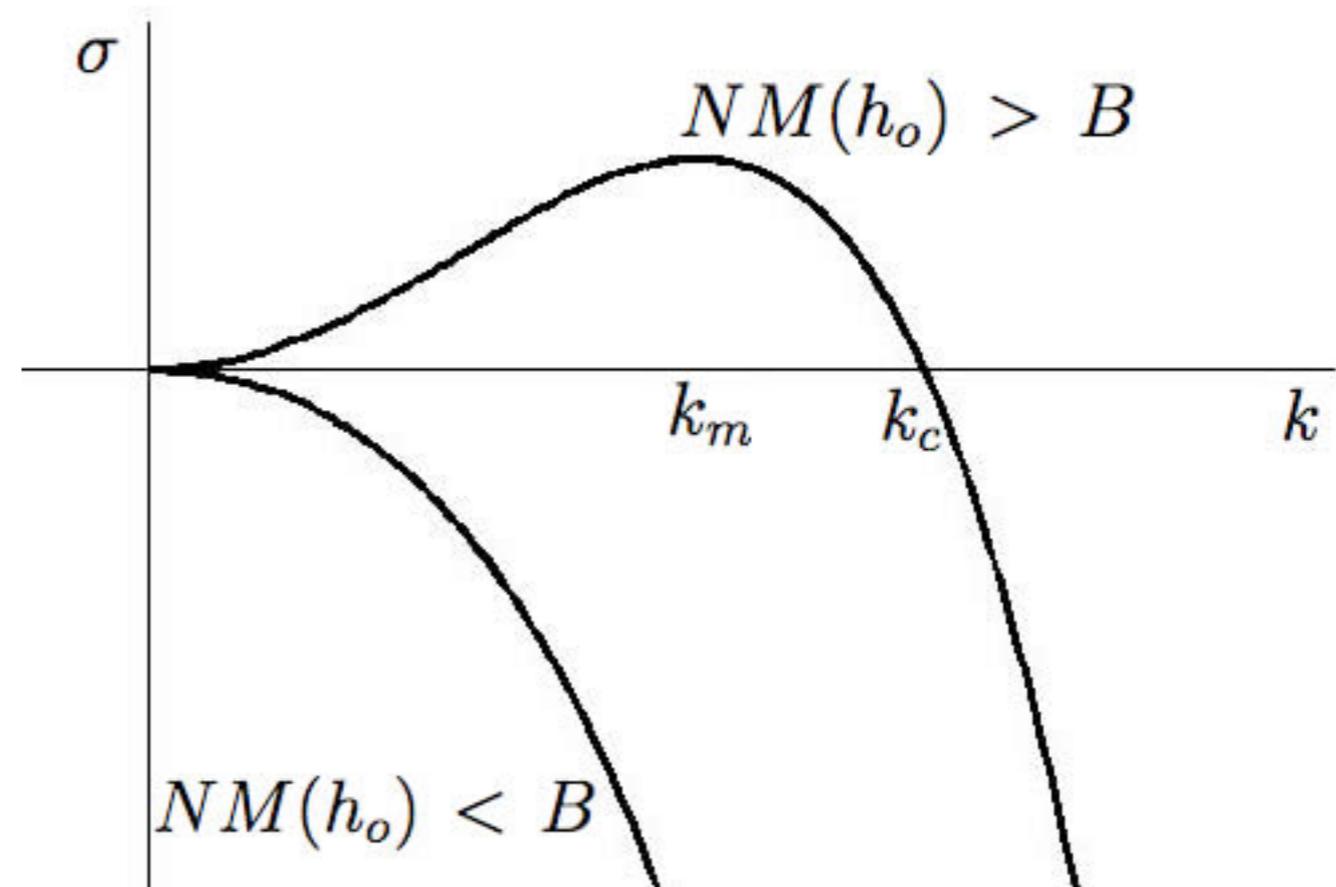
- Obtain dispersion relation by assuming solutions of the form $h_1 \propto e^{\sigma t + ikx}$

$$\sigma = -Ch_o^3 k^4 + (B - NM(h_o))h_o^3 k^2 - i3Uh_o^2 k$$

- Surface tension responsible for stabilizing system for perturbations of short wavelengths

$$k_m = \sqrt{\frac{NM(h_o) - B}{2C}}$$

$$\sigma_m = \frac{(NM(h_o) - B)^2}{4C} h_o^3$$



σ : growth-rate

k : wavenumber

k_m : fastest growing wavenumber

Traveling-wave solutions for 2D equation

$$h_t + \partial_x [h^3 (Ch_{xxx} - Bh_x) + N(m^2 - h m m') h_x] + U(h^3)_x = 0$$

$$m(h) = \frac{h^{3/2}}{\beta^{3/2}}$$

$$h \rightarrow b \text{ as } x \rightarrow -\infty \quad h \rightarrow 1 \text{ as } x \rightarrow \infty$$

$$h_x \rightarrow 0 \text{ as } x \rightarrow -\infty \quad h_x \rightarrow 0 \text{ as } x \rightarrow \infty$$

- ◉ simplify analysis by assuming $\beta \gg h$
- ◉ Neumann BCs used since far behind & far in front of the contact line, fluid thickness approximately constant

Traveling-wave solutions cont.

- ◉ Look at solution in a moving reference frame
- ◉ make a change of variables $\xi = x - Vt$ where V is the wave speed
- ◉ \therefore substituting $h_o(\xi) = h(x, t)$ and integrating

$$-Vh_o + Ch_o^3 h_{o\xi\xi\xi} - \left(B + \frac{N}{2\beta^3} \right) h_o^3 h_{o\xi} + U(h_o^3)_\xi = d$$

- ◉ Applying the BCs

$$d = -b(1 + b)$$

$$V = U(1 + b + b^2)$$

LSA for 3D evolution equation

$$h_t + \nabla \cdot [h^3 (C \nabla \nabla^2 h - B \nabla h) + N(m^2 - h m m') \nabla h] + U(h^3)_x = 0$$

- ◉ constant flux-driven case
- ◉ fluid is being injected into the film
- ◉ infinite volume
- ◉ initially, flow in transverse direction fairly stable
- ◉ apply perturbations to leading order equation

$$h(x, y, t) = h_o(\xi) + \epsilon \phi(\xi) e^{\sigma t + iky} + O(\epsilon^2)$$

LSA for 3D evolution equation cont.

$$-\sigma\phi = -Vg_\xi + C\left[k^4 h_o^3 \phi - k^2 \left((h_o^3 \phi_\xi)_\xi + h_o^3 \phi_{\xi\xi} \right) + (h_o^3 \phi_{\xi\xi\xi} + 3h_o^3 h_{o\xi\xi\xi} \phi)_\xi \right] \\ + \left(B + \frac{N}{2\beta^3} \right) \left(k^2 h_o^3 \phi - (h_o^3 \phi_\xi + 3h_o^2 h_{o\xi} \phi)_\xi \right) + 3U (h_o^2 \phi)_\xi$$

- σ, ϕ depend only on even powers of k
- looking at the solution in the limit of a small wavenumber

$$\phi = \phi_o + k^2 \phi_1 + O(k^4) \quad \sigma = \sigma_o + k^2 \sigma_1 + O(k^4)$$

- we modified the position of the contact line by the perturbation, so BCs need to be linearized accordingly giving $\phi_o(\xi) = h_{o\xi}(\xi)$

LSA for 3D evolution equation cont.

- leading-order equation

$$-\sigma_o h_{o\xi} = \left[-Vh_o + Ch_o^3 h_{o\xi\xi\xi} - \left(B + \frac{N}{2\beta^3} \right) h_o^3 h_{o\xi} + U(h_o^3)_\xi \right]_{\xi\xi}$$

- right hand side is our leading order equation, therefore $\sigma_o = 0$

LSA for 3D evolution equation cont.

- $O(k^2)$ equation

$$-\sigma_1 h_{o\xi} = -V\phi_{1\xi} + C\left(-\left(h_o^3 h_{o\xi\xi}\right)_\xi - h_o^3 h_{o\xi\xi\xi} + \left(h_o^3 \phi_{1\xi\xi\xi}\right)_\xi + 3\left(h_o^2 h_{o\xi\xi\xi} \phi_1\right)_\xi\right) \\ + \left(B + \frac{N}{2\beta^3}\right)\left(h_o^3 h_{o\xi} - \left(h_o^3 \phi_1\right)_{\xi\xi}\right) + 3U\left(h_{o\xi} \phi_1\right)_\xi$$

- We integrate and apply the BCs
- Finally,

$$\sigma \approx \frac{k^2}{1-b} \int_{-\infty}^0 C h_o^3 h_{o\xi\xi\xi} - \left(B + \frac{N}{2\beta^3}\right) h_o^3 h_{o\xi} d\xi$$