

Physics 121 - Electricity and Magnetism

Lecture 06 - Capacitance

Y&F Chapter 24 Sec. 1 - 6

- **Overview**
- **Definition of Capacitance**
- **Calculating the Capacitance**
- **Parallel Plate Capacitor**
- **Spherical and Cylindrical Capacitors**
- **Capacitors in Parallel and Series**
- **Energy Stored in an Electric Field**
- **Atomic Physics View of Dielectrics**
- **Electric Dipole in an Electric Field**
- **Capacitors with a Dielectric**
- **Dielectrics and Gauss Law**
- **Summary**

What Capacitance Measures

How much charge does an arrangement of conductors hold when a given voltage is applied?

- The charge needed depends on a geometrical factor called capacitance.

$$Q = C\Delta V$$

Example:

- Two conducting spheres: Radii R_1 and $R_2 = 2R_1$. Different charges Q_1 and Q_2 .
- Spheres touch and come to the same potential ΔV ,
- Apply point charge potential formula, $V(\text{infinity}) = 0$

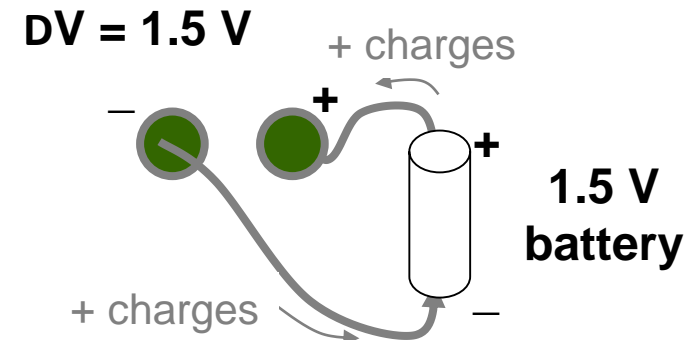
$$\Delta V = \frac{1}{4\pi\epsilon_0} \frac{Q_1}{R_1} \equiv \frac{Q_1}{C_1} \quad \text{and also} \quad \Delta V = \frac{1}{4\pi\epsilon_0} \frac{Q_2}{R_2} \equiv \frac{Q_2}{C_2} \quad \rightarrow \quad \frac{Q_1}{Q_2} \equiv \frac{C_1}{C_2} = \frac{R_1}{R_2} = \frac{1}{2}$$

Capacitance of a single isolated sphere:

$$C = 4\pi\epsilon_0 R$$

Example: A primitive capacitor

- The right ball's potential is the same as the + side of the battery. Similarly for the - ball.
- How much charge flows onto each ball to produce a potential difference of 1.5 V?
- The answer depends on the capacitance.



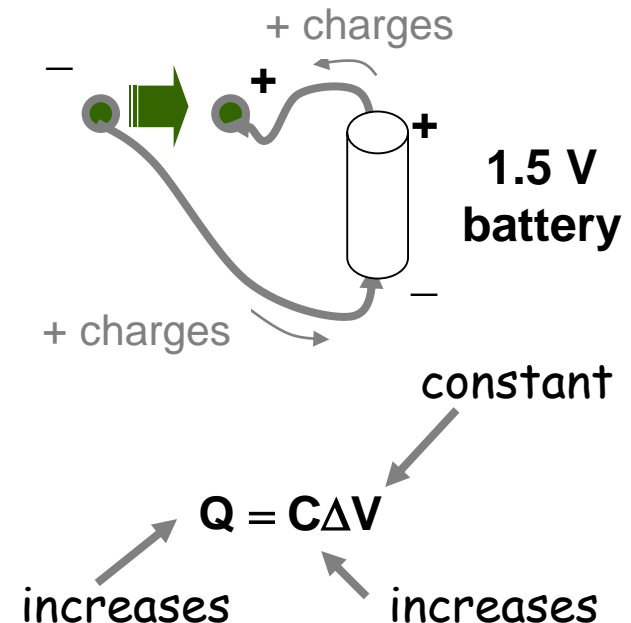
Definition of CAPACITANCE :

$$C \equiv \frac{Q}{\Delta V} \quad \text{or} \quad Q \equiv C \Delta V \quad \left[\frac{\text{Coulombs}}{\text{Volt}} \right]$$

- Measures the charge needed per volt of potential difference
- Does not depend on applied ΔV or charge Q . Always positive.
- Depends on geometry (and on dielectric materials) only
- Units: 1 FARAD = 1 Coulomb / Volt. - Farads are very large
 1 mF = 10^{-6} F. 1 pF = 1 pico-Farad = 10^{-12} F = 10^{-6} μ F = 1 $\mu\mu$ F

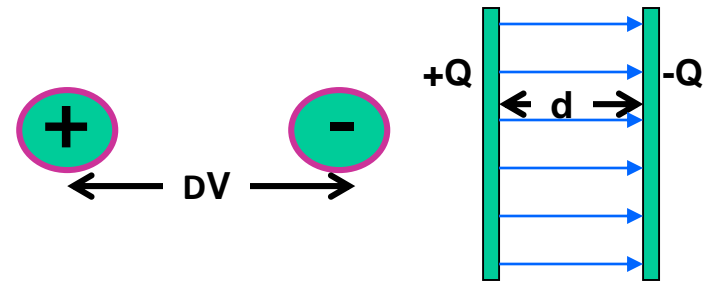
Example - Capacitance depends on geometry

- Move the balls at the ends of the wires closer together while still connected to the battery
- The potential difference ΔV cannot change.
- But: $\Delta V = -\int \vec{E} \cdot d\vec{s} \approx \vec{E}_{av} \circ \Delta \vec{s}$
- The distance D s between the balls decreased so the E field had to increase as did the stored energy.
- Charge flowed from the battery to the balls to increase E .
- The two balls now hold more charge for the same potential difference: i.e. the capacitance increased.



Capacitors are charge storage devices

- Two conductors not in electrical contact
- Electrically neutral before & after being charged
 - $Q_{\text{enc}} = Q_{\text{net}} = 0$
- Current can flow from + plate to - plate if there is a conducting path (complete circuit)
- Capacitors store charge and potential energy
 - memory bits - radio circuits - power supplies
- Common type: “parallel plate”, sometimes tubular



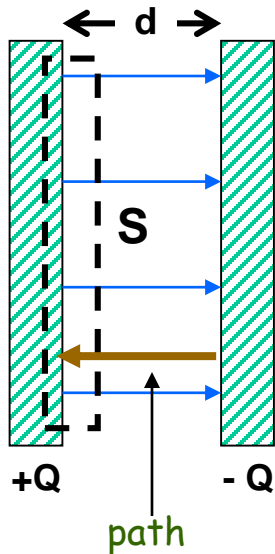
Method for calculating capacitance from geometry:

- Assume two conducting plates (equipotentials) with equal and opposite charges +Q and -Q
- Possibly use Gauss' Law to find E between the plates
- Calculate ΔV between plates using a convenient path
- Capacitance $C = Q/\Delta V$
- Certain materials (“dielectrics”) can reduce the E field between plates by “polarizing” - capacitance increases

$$\Phi_E = \frac{q_{\text{enc}}}{\epsilon_0} = \oint_S \vec{E} \cdot d\vec{A}$$

$$\Delta V_{fi} = -\int_f^i \vec{E} \cdot d\vec{s}$$

EXAMPLE: CALCULATE C for a PARALLEL PLATE CAPACITOR



Find E between plates

- A = plate area. Treat plates as infinite sheets
- $|s^+| = |-s^-| = s = Q / A =$ uniform surface charge density
- E is uniform between the plates ($d \ll$ plate size)
- Use Gaussian surface S (one plate). Flux through ends and attached conductors is zero. Total flux is EA
- $Q_{\text{enc}} = \sigma A = e_0 f = \epsilon_0 EA$

$$\therefore E = \sigma / \epsilon_0 \quad \text{i.e.,} \quad E = Q / \epsilon_0 A$$

(infinite conducting sheet)

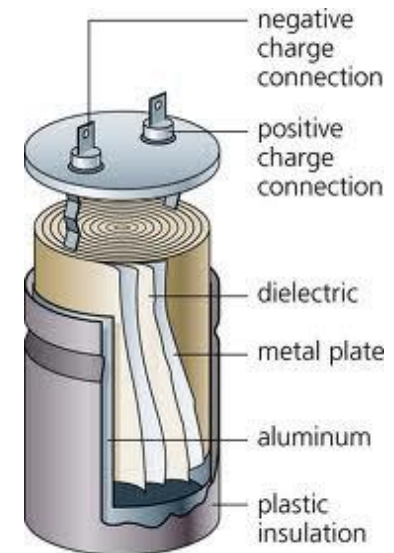
Find potential difference DV:

- Choose $V = 0$ on negative plate (grounded)
- Choose path from - plate to + plate, opposite to E field

$$\Delta V_{\text{fi}} = - \int_{\text{path}} \vec{E} \cdot d\vec{s} = (-)(-)Ed = + \frac{\sigma}{\epsilon_0} d = \frac{Qd}{\epsilon_0 A} \equiv \frac{Q}{C}$$

$$\therefore C \equiv \frac{Q}{\Delta V} = \frac{\epsilon_0 A}{d}$$

- C DEPENDS ONLY ON GEOMETRY
- $C \rightarrow$ infinity as plate separation $d \rightarrow 0$
- C directly proportional to plate area A



- Other formulas for other geometries

EX 24.03: FIND C for a SPHERICAL CAPACITOR

- 2 concentric spherical, conducting shells, radii a & b
- Charges are $+q$ (inner sphere), $-q$ (outer sphere)
- All charge on the outer sphere is on its inner surface (by Gauss's Law)
- Choose Gaussian surface S as shown and find field using Gauss's Law:

$$\epsilon_0 \int_S \vec{E} \cdot d\vec{A} = q \quad q = \epsilon_0 \mathbf{E}A = \epsilon_0 \mathbf{E}(4\pi r^2)$$

- As before: $\mathbf{E} = q/(4\pi\epsilon_0 r^2)$
- To find potential difference use outward radial integration path from $r = a$ to $r = b$.

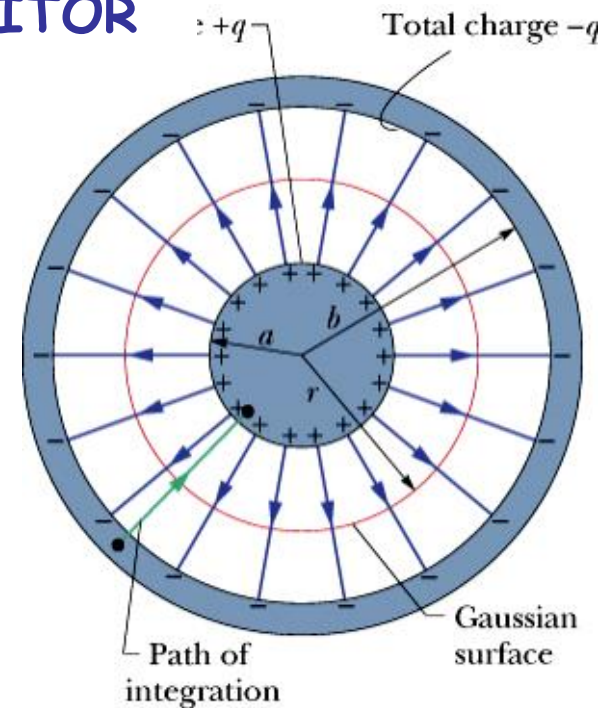
$$\Delta V = V_b - V_a = - \int_{r=a}^{r=b} \vec{E} \cdot d\vec{s} = \frac{-q}{4\pi\epsilon_0} \int_{r=a}^{r=b} \frac{dr}{r^2} = \frac{-q}{4\pi\epsilon_0} \left. \frac{(-1)}{r} \right|_a^b$$

$$\Delta V = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{b} - \frac{1}{a} \right) = \frac{q}{4\pi\epsilon_0} \left(\frac{a-b}{ba} \right) \quad \text{Negative For } b > a \quad \mathbf{V_b < V_a}$$

$$\therefore \mathbf{C} \equiv \frac{q}{|\Delta V|} = \frac{4\pi\epsilon_0 ab}{b-a}$$

Let $b \rightarrow$ infinity. Then $a/b \rightarrow 0$ and result becomes the earlier formula for the isolated sphere:

$$\mathbf{C} \rightarrow \frac{4\pi\epsilon_0 ab}{b} = 4\pi\epsilon_0 a$$



EX 24.04: Find C for a **CYLINDRICAL CAPACITOR**

- 2 concentric, long cylindrical conductors
- Radii a & b and length $L \gg b \Rightarrow$ neglect end effects
- Charges are $+q$ (inner) and $-q$ (outer), λ is uniform
- All charge on the outer conductor is on its inner surface (by Gauss's Law)
- Choose Gaussian surface S between plates and find field at radius r .
- E is perpendicular to endcaps \Rightarrow zero flux contribution

$$\epsilon_0 \Phi_{\text{cyl}} = \epsilon_0 \int_S \vec{E} \cdot d\vec{A} = q \quad q = \epsilon_0 EA = \epsilon_0 E(2\pi rL)$$

- So: $E = q/(2\pi\epsilon_0 rL) = \lambda/(2\pi\epsilon_0 r)$

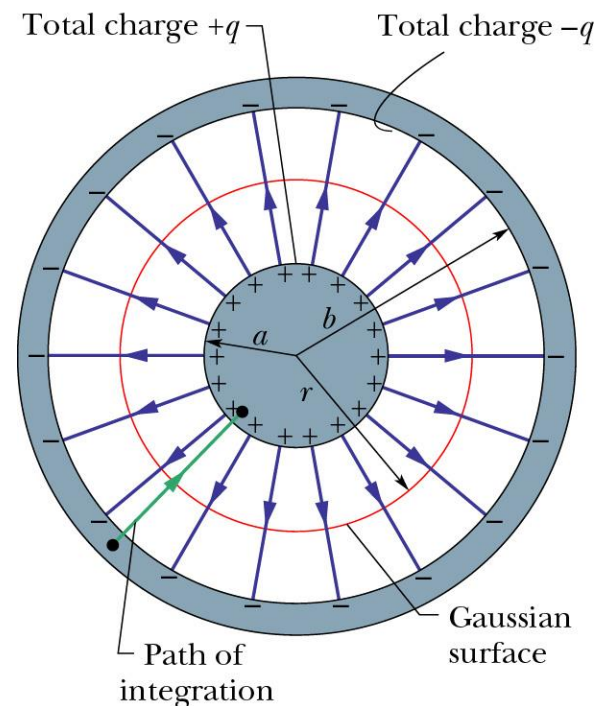
- To find potential difference use outward radial integration path from $r = b$ to $r = a$.

$$\Delta V = V_b - V_a = - \int_{r=a}^{r=b} \vec{E} \cdot d\vec{s} = \frac{-q}{2\pi\epsilon_0 L} \int_{r=a}^{r=b} \frac{dr}{r} = \frac{-q}{2\pi\epsilon_0 L} \ln(r) \Big|_a^b = \frac{-q}{2\pi\epsilon_0 L} \ln(b/a)$$

$$C = q/\Delta V = 2\pi\epsilon_0 \frac{L}{\ln(b/a)}$$

$$C \rightarrow 0 \text{ as } b/a \rightarrow \text{inf}$$

$$C \rightarrow \text{inf as } b/a \rightarrow 1$$



$$V_b < V_a \text{ For } b > a$$

C depends only on
geometrical parameters

Examples of Capacitance Formulas

- **Capacitance for isolated Sphere**

$$C = 4\pi\epsilon_0 R$$

- **Parallel Plate Capacitor**

$$C = \frac{\epsilon_0 A}{d}$$

- **Concentric Cylinders Capacitor**

$$C = 2\pi\epsilon_0 \frac{L}{\ln(b/a)}$$

- **Concentric Spheres Capacitor**

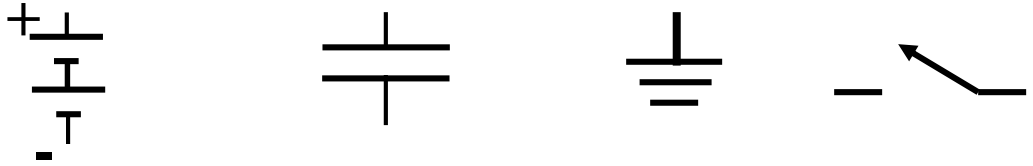
$$C = 4\pi\epsilon_0 \frac{ab}{b-a}$$

- **Units: F (Farad) = C²/Nm = C/ Volt = $\epsilon_0 \times \text{length}$**
 - **named after Michael Faraday. [note: $\epsilon_0 = 8.85 \text{ pF/m}$]**

All of these formulas depend only on geometrical factors

Capacitors in circuits

CIRCUIT SYMBOLS:



CIRCUIT DEFINITIONS:

Current $\equiv i \equiv dq/dt$
= rate of + charge flow past a point in the circuit

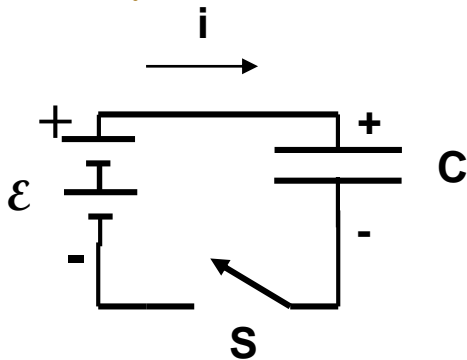
Open Circuit: NO closed path. No current. Conductors are equi-potentials

Closed Circuit: There is/are completed paths through which current can flow.

Loop Rule: Potential is a conservative field

→ Potential CHANGE around ANY closed path = 0

Example: CHARGING A CAPACITOR



- Current flows when switch is CLOSED, completing circuit
- Battery (EMF) maintains DV (= EMF \mathcal{E}), and supplies energy by moving free + charges from - to + terminal, internal to battery

Convention: i flows from + to - outside of battery

When switch closes, current (charge) flows until DV across capacitor equals battery voltage \mathcal{E} .

Then current stops as E field in wire $\rightarrow 0$

DEFINITION: EQUIVALENT CAPACITANCE

- Capacitors can be connected in series, parallel, or more complex combinations
- The “equivalent capacitance” is the capacitance of a SINGLE capacitor that would have the same capacitance as the combination.
- The equivalent capacitance can replace the original combination in analysis.

Parallel capacitors - Equivalent capacitance

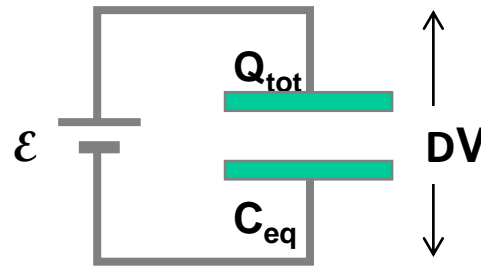
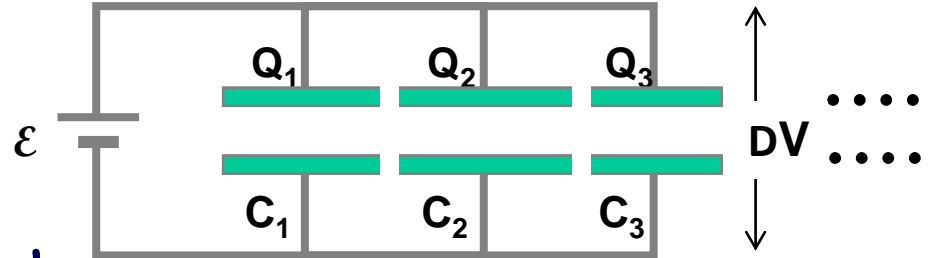
The actual parallel circuit...

$$Q_i = C_i \Delta V$$

ΔV is the same for each branch

...and the equivalent circuit:

$$Q_{\text{tot}} \equiv C_{\text{eq}} \Delta V$$



The parallel capacitors are just like a single capacitor with larger plates so....

$$Q_{\text{tot}} = \sum Q_i \quad (\text{parallel})$$

Charges on parallel capacitors add

$$\therefore Q_{\text{tot}} = C_1 \Delta V + C_2 \Delta V + C_3 \Delta V + \dots = (C_1 + C_2 + C_3 + \dots) \Delta V$$

Parallel capacitances add directly

$$C_{\text{eq}} = \sum C_i \quad (\text{parallel})$$

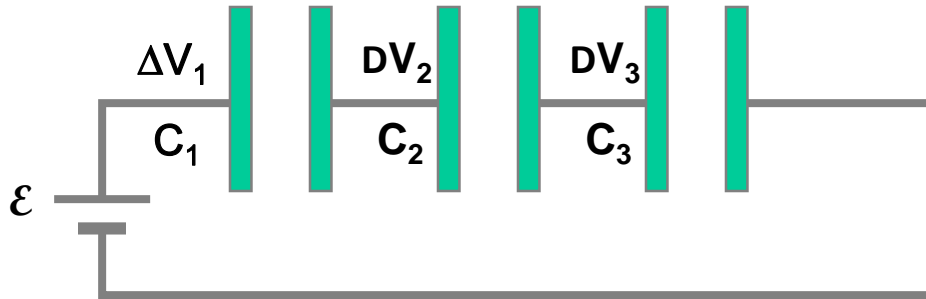
Question: Why is ΔV the same for all elements in parallel?

Answer: Potential is conservative field, for ANY closed loop around circuit:

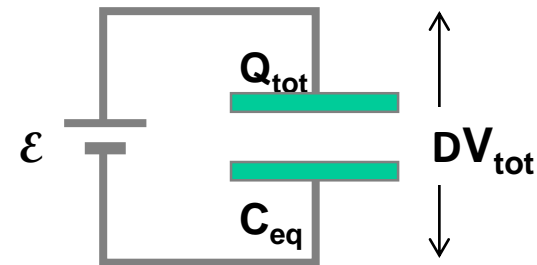
$$\sum \Delta V_i = 0 \quad (\text{Kirchoff Loop Rule})$$

Series capacitors - equivalent capacitance

The actual series circuit...



The equivalent circuit...



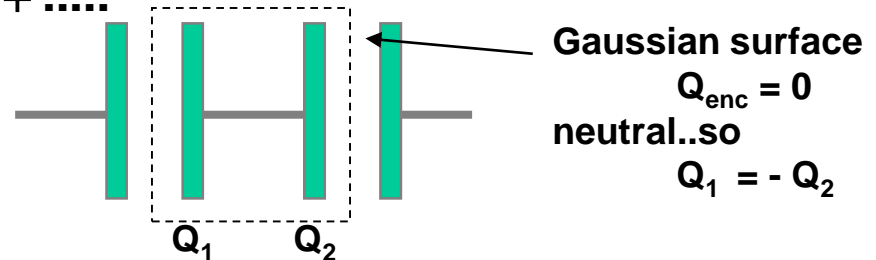
ΔV_i are NOT necessarily the same for each capacitor in series

$$\Delta V_{\text{tot}} = \sum \Delta V_i = \Delta V_1 + \Delta V_2 + \Delta V_3 + \dots$$

But... charges on series capacitors are all equal - here's why.....

so $\Delta V_i = Q/C_i$ same Q

$$\therefore \Delta V_{\text{tot}} = Q/C_1 + Q/C_2 + Q/C_3 + \dots = Q \sum 1/C_i \equiv Q/C_{\text{eq}}$$



Reciprocals of series capacitances add

$$\frac{1}{C_{\text{eq}}} = \sum \frac{1}{C_i} \quad (\text{series})$$

For two capacitors in series:

$$\frac{1}{C_{\text{eq}}} = \frac{1}{C_1} + \frac{1}{C_2} = \frac{C_2 + C_1}{C_1 C_2} \Rightarrow C_{\text{eq}} = \frac{C_2 C_1}{C_1 + C_2}$$

Example 1: A $33\mu\text{F}$ and a $47\mu\text{F}$ capacitor are connected in parallel
Find the equivalent capacitance

Solution: $C_{\text{para}} = C_1 + C_2 = 80\mu\text{F}$

Example 2: Same two capacitors as above, but now in series connection

Solution: $C_{\text{ser}} = \frac{C_1 C_2}{C_1 + C_2} = \frac{33 \times 47}{33 + 47} = 19.4\mu\text{F}$

Example 3: A pair of capacitors is connected as shown



- $C_1 = 10\mu\text{F}$, charged initially to $100\text{V} = V_i$
- $C_2 = 20\mu\text{F}$, uncharged initially

Close switches. Find final potentials across C_1 & C_2 .

Solution:

- C 's are in parallel \rightarrow Same potential V_f for each
- Total initial charge: $Q_{\text{tot}} = Q_{1i} = C_1 V_i = 10^{-3}\text{ C}$.
- Charge is conserved - it redistributes on both C_1 & C_2

$$C_{\text{eq}} = Q_{\text{tot}} / V_f = C_1 + C_2 \Rightarrow V_f = \frac{10^{-3}}{30 \times 10^{-6}} = 33\text{ V}.$$

- Final charge on each:

$$Q_{1f} = C_1 V_f = 3.3 \times 10^{-4}\text{ C}. \quad Q_{2f} = C_2 V_f = 6.7 \times 10^{-4}\text{ C}.$$

Three Capacitors in Series

6-2: The equivalent capacitance for two capacitors in series is:

$$C_{eq} = \frac{1}{\frac{1}{C_1} + \frac{1}{C_2}} = \frac{C_1 C_2}{C_1 + C_2}$$



Which of the following is the equivalent capacitance formula for three capacitors in series?

A.
$$C_{eq} = \frac{C_1 C_2 C_3}{C_1 + C_2 + C_3}$$

B.
$$C_{eq} = \frac{C_1 C_2 + C_2 C_3 + C_1 C_3}{C_1 + C_2 + C_3}$$

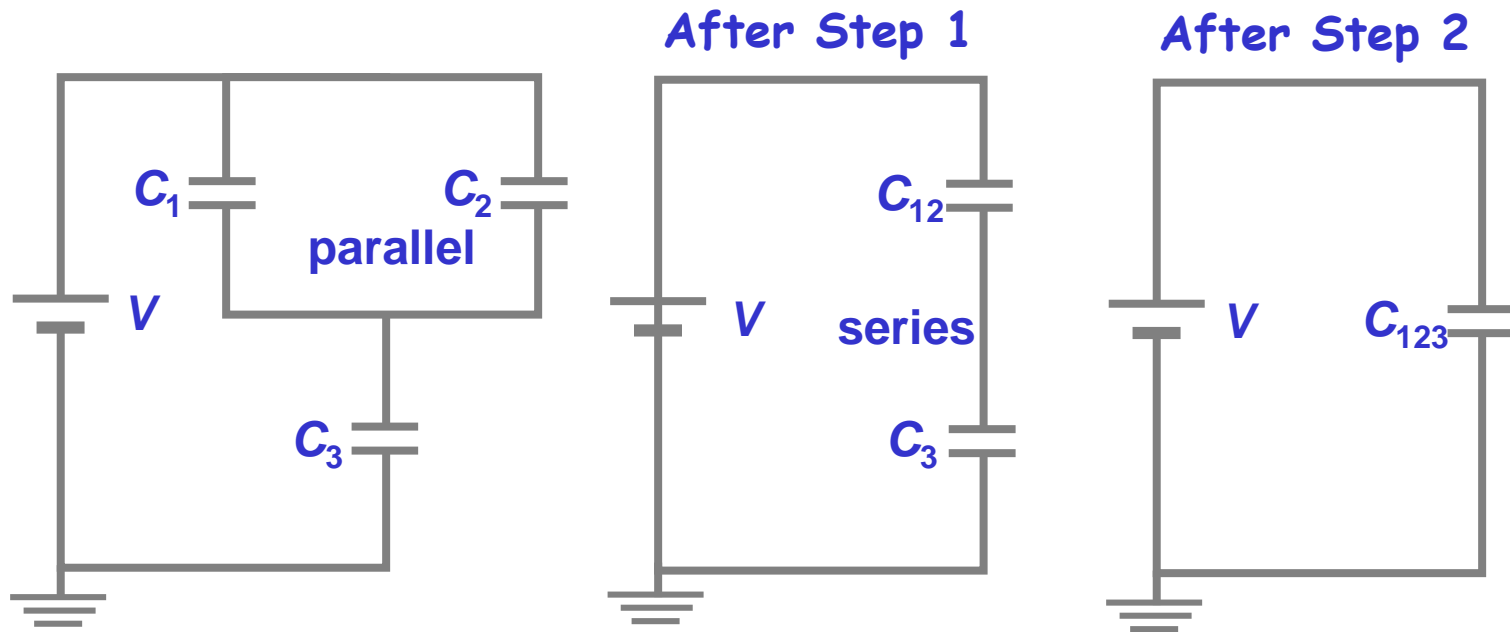
C.
$$C_{eq} = \frac{C_1 C_2 + C_2 C_3 + C_3 C_1}{C_1 C_2 C_3}$$

D.
$$C_{eq} = \frac{C_1 + C_2 + C_3}{C_1 C_2 C_3}$$

E.
$$C_{eq} = \frac{C_1 C_2 C_3}{C_1 C_2 + C_2 C_3 + C_3 C_1}$$

Apply formula for C_{eq} twice

Example: Reduce circuit to find $C_{eq} = C_{123}$ for mixed series-parallel capacitors



$$C_{12} = C_1 + C_2$$

$$\frac{1}{C_{123}} = \frac{1}{C_{12}} + \frac{1}{C_3}$$

$$C_{123} = \frac{C_{12}C_3}{C_{12} + C_3}$$

Values:

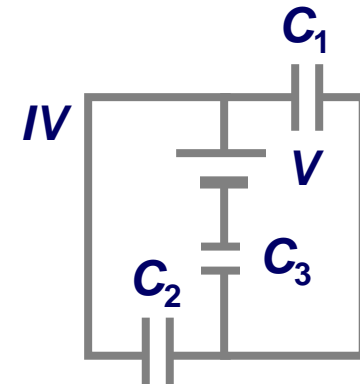
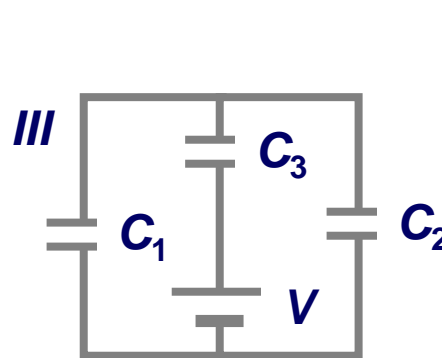
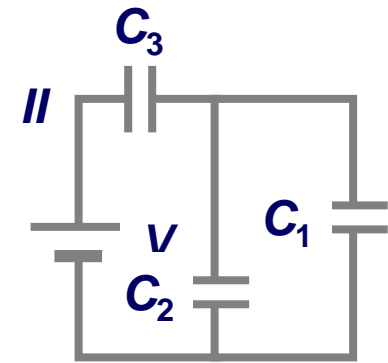
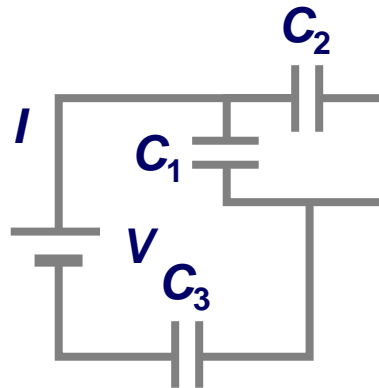
$$C_1 = 12.0 \mu\text{F}, C_2 = 5.3 \mu\text{F}, C_3 = 4.5 \mu\text{F}$$

$$C_{123} = (12 + 5.3) 4.5 / (12 + 5.3 + 4.5) \mu\text{F} = 3.57 \mu\text{F}$$

Series or Parallel?

6-3: In the circuits below, which ones show capacitors 1 and 2 connected in series?

- A. I, II, III
- B. I, III
- C. II, IV
- D. III, IV
- E. None



Energy Stored in a Capacitor

When charge flows in the sketch, energy stored in the battery is depleted. Where does it go?

- Charge distributions have potential energy. Charges that are separated in a neutral body store energy.
- The electric potential is defined to be

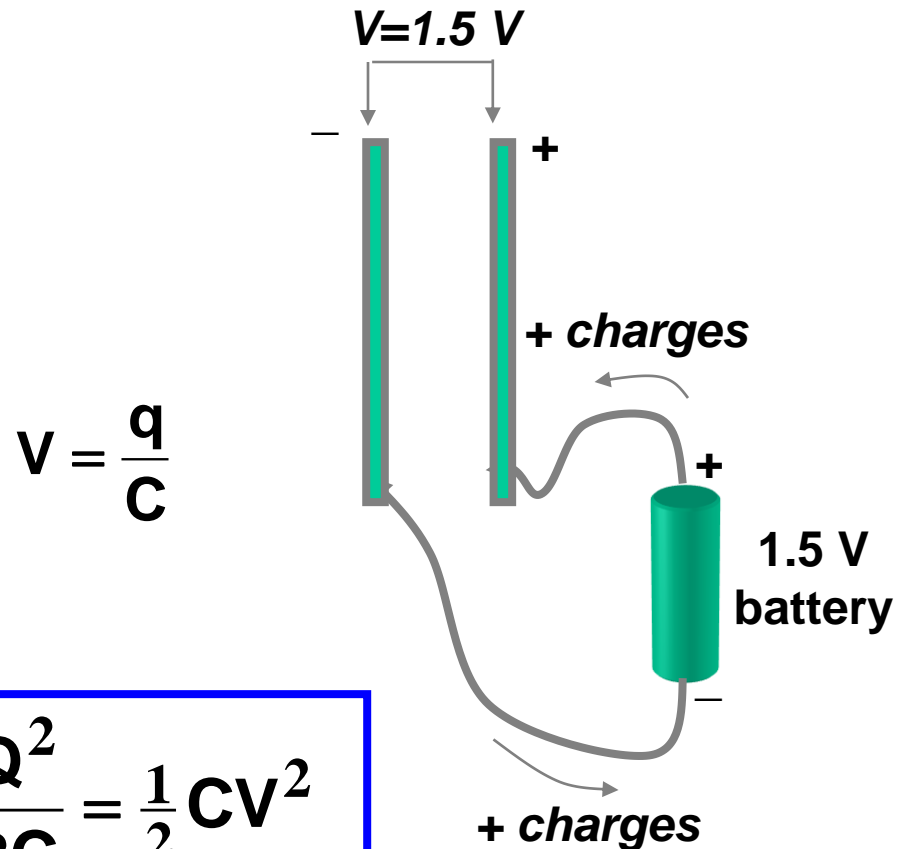
$$V = U/q, \quad U = qV$$

- A small element of charge dq on each plate of a capacitor stores potential energy:

$$dU = V dq$$

- The energy stored by charging a capacitor from charge 0 to Q is the integral:

$$U = \int_0^Q dU = \frac{1}{C} \int_0^Q q' dq' = \frac{Q^2}{2C} = \frac{1}{2} CV^2$$



Capacitors Store Energy in the Electrostatic Field

- The total energy in a parallel plate capacitor is

$$U = \frac{1}{2} CV^2 = \frac{\epsilon_0 A}{2d} V^2$$

- The volume of space filled by the electric field in the capacitor is $= Ad$, so the *energy density* u is

$$u \equiv \frac{U}{\text{vol}} = \frac{\epsilon_0 A}{2dAd} V^2 = \frac{1}{2} \epsilon_0 \left(\frac{V}{d} \right)^2$$

- But for a parallel plate capacitor,

$$V = -\int \vec{E} \cdot d\vec{s} = Ed$$

- so

$$u = \frac{1}{2} \epsilon_0 E^2$$

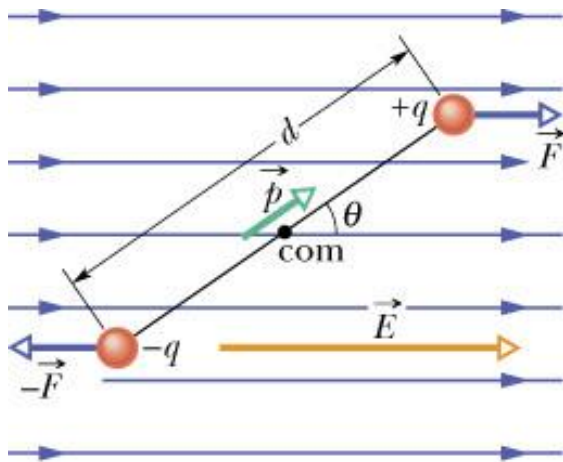
Energy is stored in
the electric field



Model for a Molecule that can Polarize

A dipole in a uniform external field..

..feels torque, stores electrostatic potential energy

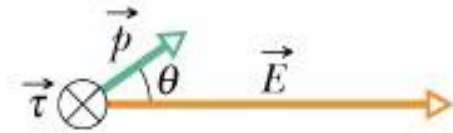


See Lecture 3

$$\vec{p} \equiv q\vec{d}$$

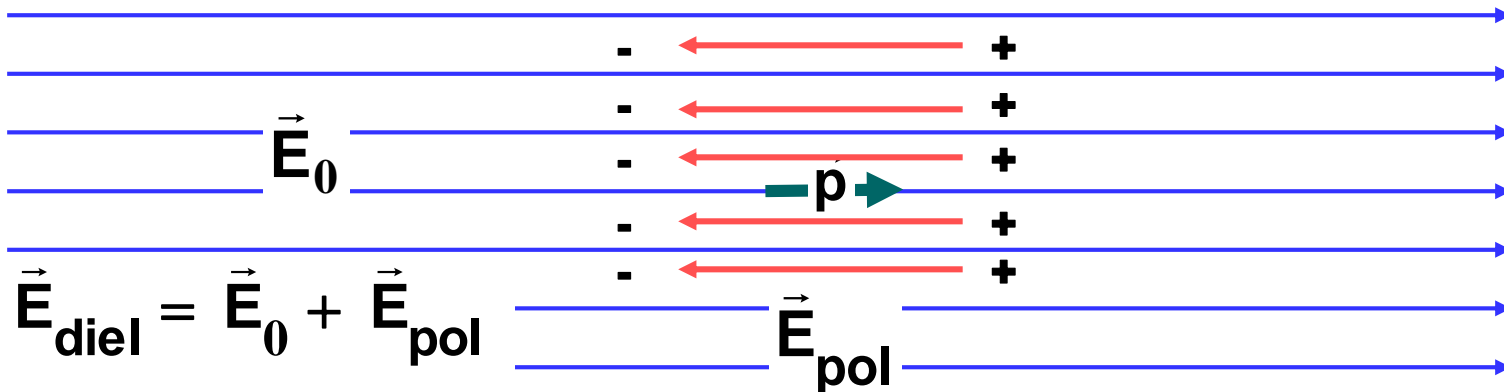
$$\vec{\tau} = \vec{p} \times \vec{E}$$

- |torque| = 0 at $\alpha = 0$ or $\alpha = \pi$
- |torque| = pE at $\alpha = \pm \pi/2$
- RESTORING TORQUE: $t(-q) = t(+q)$



$$U_E = -\vec{p} \cdot \vec{E}$$

Polarization: An external field aligns dipoles in a material, causing polarization that reduces the field



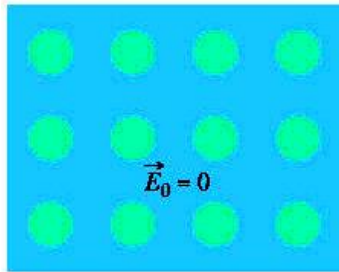
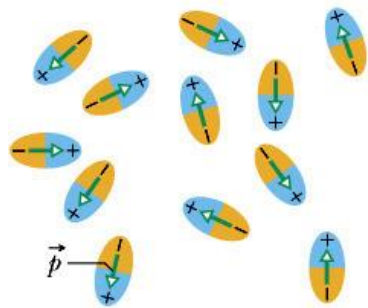
Dielectric materials in capacitors

- Insulators *POLARIZE* when an external electric field is applied
- The NET field inside the material is reduced.

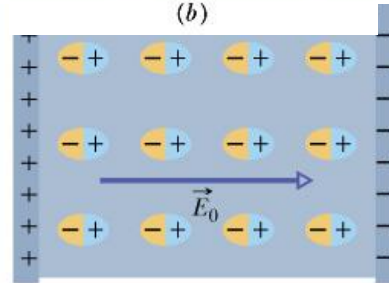
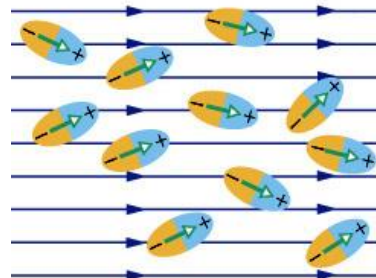
MOLECULAR VIEW

Polar Material permanent dipole
Non-polar Material induced dipole

NO EXTERNAL FIELD



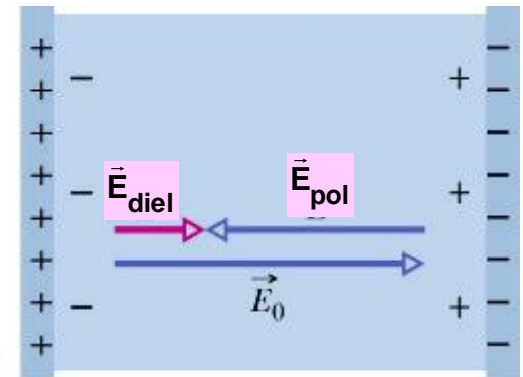
WITH EXTERNAL FIELD E_0



$E_0 = E_{vac}$ is field due to free charge
Response to E_{vac} is the polarization field E_{pol}

Actual weakened net field inside is

$$E_{diel} = E_0 - E_{pol} = E_{net}$$



Polarization surface charge density reduces free surface charge density

$$\sigma_{net} = s_{free} - s_{pol}$$

$$(s_{free} = s_{ext} = s_{vac})$$

Dielectrics increase capacitance
For a given DV, more movable charge s_{free} is needed

Dielectric constant $\equiv K \equiv \frac{C_{dielectric}}{C_{vacuum}} \geq 1$

Inside conductors, polarization reduces E_{net} to zero

Representing Dielectrics

- ϵ_0 is the free space permittivity.
- All materials (water, paper, plastic, air) polarize to some extent and have different permittivities $\epsilon = \kappa\epsilon_0$
- κ is the *dielectric constant* - a dimensionless number.
- Wherever you see ϵ_0 for a vacuum, you can substitute $\kappa\epsilon_0$ when considering dielectric materials.
- For example, the capacitance of a parallel plate capacitor increases when the space is filled with a dielectric:

$$C_{\text{diel}} = \frac{\kappa\epsilon_0 A}{d} = \kappa C_{\text{vac}}$$

- A dielectric weakens the field, compared to what it would be for a vacuum

$$\vec{E}_{\text{diel}} = \vec{E}_{\text{vac}} / \kappa \equiv \vec{D} / \epsilon_0 \kappa$$

TABLE 26.1**Approximate Dielectric Constants and Dielectric Strengths of Various Materials at Room Temperature**

Material	Dielectric Constant κ	Dielectric Strength^a (10^6 V/m)
Air (dry)	1.000 59	3
Bakelite	4.9	24
Fused quartz	3.78	8
Mylar	3.2	7
Neoprene rubber	6.7	12
Nylon	3.4	14
Paper	3.7	16
Paraffin-impregnated paper	3.5	11
Polystyrene	2.56	24
Polyvinyl chloride	3.4	40
Porcelain	6	12
Pyrex glass	5.6	14
Silicone oil	2.5	15
Strontium titanate	233	8
Teflon	2.1	60
Vacuum	1.000 00	—
Water	80	—

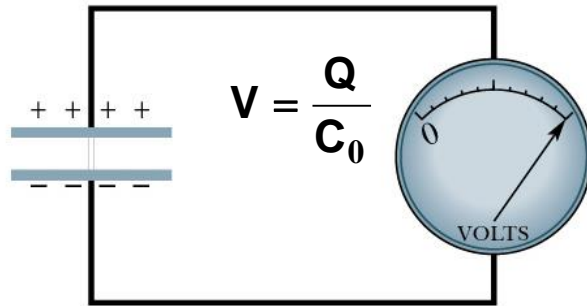
^a The dielectric strength equals the maximum electric field that can exist in a dielectric without electrical breakdown. These values depend strongly on the presence of impurities and flaws in the materials.

What happens as you insert a dielectric?

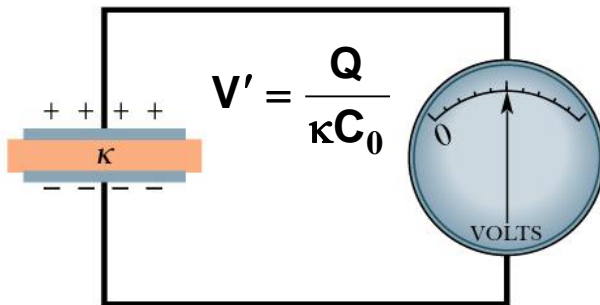
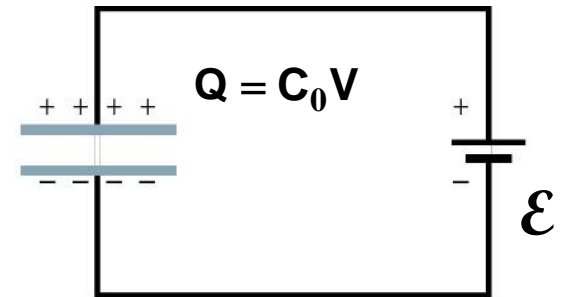
Initially, charge capacitor C_0 to voltage V , charge Q , field E_{net} .

- With battery **detached** insert dielectric
- Q remains constant, E_{net} is reduced
- Voltage (fixed Q) drops to V' .
- Dielectric reduced E_{net} and V .

- With battery **attached**, insert dielectric.
- E_{net} and V are momentarily reduced but battery maintains voltage \mathcal{E}
- Charge flows to the capacitor as dielectric is inserted until V and E_{net} are back to original values.

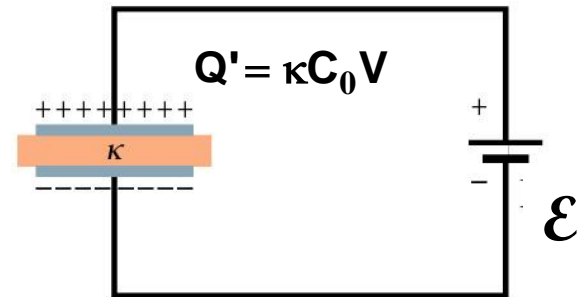


OR



$Q = \text{a constant}$

$$V = -\int \vec{E} \cdot d\vec{s} = Ed$$



$V = \text{a constant}$

$Q' > Q$

Gauss' Law with a dielectric

$$\epsilon_0 \int_S \vec{E}_{\text{diel}} \circ d\vec{A} = q_{\text{free}} - q_{\text{pol}} = q_{\text{net}}$$

OPTIONAL
TOPIC

Alternatively:

$$\int_S \underbrace{\epsilon_0 \vec{E}_{\text{vac}}}_{\text{field not counting polarization} = \epsilon_0 \vec{E}_{\text{vac}}} \circ d\vec{A} = q_{\text{free}} = \int_S \epsilon_0 K \vec{E}_{\text{diel}} \circ d\vec{A} \equiv \int_S \vec{D} \circ d\vec{A}$$

free charge on plates

The “Electric Displacement” D measures field that would be present due to the “free” charge only, i.e. without polarization field from dielectric

$$\vec{E}_{\text{vac}} = K \vec{E}_{\text{diel}} \quad \vec{D} \equiv \epsilon_0 \vec{E}_{\text{vac}} = \epsilon_0 K \vec{E}_{\text{diel}} \equiv \epsilon \vec{E}_{\text{diel}}$$

- K could vary over Gaussian surface S . Usually it is constant and factors
- Flux is still measured using field **without** dielectric: $E_{\text{vac}} = K E_{\text{diel}} = D/\epsilon_0$

$$d\Phi = \vec{E}_{\text{vac}} \circ d\vec{A} = K \vec{E}_{\text{diel}} \circ d\vec{A}$$

- Only the **free charges** q_{free} (excluding polarization) are counted as q_{enc} in the above. Using K on the left compensates for the polarization.
- When applying the above include only q_{free} .
Ignore polarization charges inside the Gaussian surface

CHAPTER 24

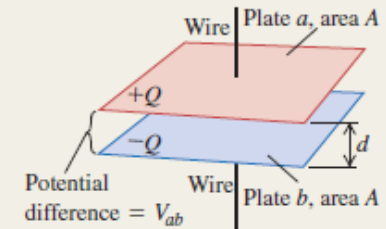
SUMMARY

Capacitors and capacitance: A capacitor is any pair of conductors separated by an insulating material. When the capacitor is charged, there are charges of equal magnitude Q and opposite sign on the two conductors, and the potential V_{ab} of the positively charged conductor with respect to the negatively charged conductor is proportional to Q . The capacitance C is defined as the ratio of Q to V_{ab} . The SI unit of capacitance is the farad (F): $1 \text{ F} = 1 \text{ C/V}$.

A parallel-plate capacitor consists of two parallel conducting plates, each with area A , separated by a distance d . If they are separated by vacuum, the capacitance depends only on A and d . For other geometries, the capacitance can be found by using the definition $C = Q/V_{ab}$. (See Examples 24.1–24.4.)

$$C = \frac{Q}{V_{ab}} \quad (24.1)$$

$$C = \frac{Q}{V_{ab}} = \epsilon_0 \frac{A}{d} \quad (24.2)$$



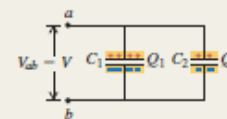
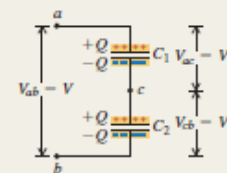
Capacitors in series and parallel: When capacitors with capacitances C_1, C_2, C_3, \dots are connected in series, the reciprocal of the equivalent capacitance C_{eq} equals the sum of the reciprocals of the individual capacitances. When capacitors are connected in parallel, the equivalent capacitance C_{eq} equals the sum of the individual capacitances. (See Examples 24.5 and 24.6.)

$$\frac{1}{C_{\text{eq}}} = \frac{1}{C_1} + \frac{1}{C_2} + \frac{1}{C_3} + \dots \quad (24.5)$$

(capacitors in series)

$$C_{\text{eq}} = C_1 + C_2 + C_3 + \dots \quad (24.7)$$

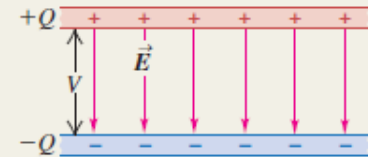
(capacitors in parallel)



Energy in a capacitor: The energy U required to charge a capacitor C to a potential difference V and a charge Q is equal to the energy stored in the capacitor. This energy can be thought of as residing in the electric field between the conductors; the energy density u (energy per unit volume) is proportional to the square of the electric-field magnitude. (See Examples 24.7–24.9.)

$$U = \frac{Q^2}{2C} = \frac{1}{2}CV^2 = \frac{1}{2}QV \quad (24.9)$$

$$u = \frac{1}{2}\epsilon_0 E^2 \quad (24.11)$$



Dielectrics: When the space between the conductors is filled with a dielectric material, the capacitance increases by a factor K , called the dielectric constant of the material. The quantity $\epsilon = K\epsilon_0$ is called the permittivity of the dielectric. For a fixed amount of charge on the capacitor plates, induced charges on the surface of the dielectric decrease the electric field and potential difference between the plates by the same factor K . The surface charge results from polarization, a microscopic rearrangement of charge in the dielectric. (See Example 24.10.)

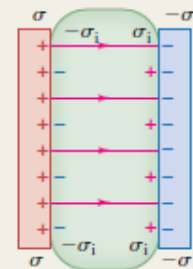
$$C = KC_0 = K\epsilon_0 \frac{A}{d} = \epsilon \frac{A}{d} \quad (24.19)$$

(parallel-plate capacitor filled with dielectric)

$$u = \frac{1}{2}K\epsilon_0 E^2 = \frac{1}{2}\epsilon E^2 \quad (24.20)$$

$$\oint K\vec{E} \cdot d\vec{A} = \frac{Q_{\text{encl-free}}}{\epsilon_0} \quad (24.23)$$

Dielectric between plates



Under sufficiently strong fields, dielectrics become conductors, a situation called dielectric breakdown. The maximum field that a material can withstand without breakdown is called its dielectric strength.

In a dielectric, the expression for the energy density is the same as in vacuum but with ϵ_0 replaced by $\epsilon = K\epsilon_0$. (See Example 24.11.)

Gauss's law in a dielectric has almost the same form as in vacuum, with two key differences: \vec{E} is replaced by $K\vec{E}$ and Q_{encl} is replaced by $Q_{\text{encl-free}}$, which includes only the free charge (not bound charge) enclosed by the Gaussian surface. (See Example 24.12.)