## Physics 121 - Electricity and Magnetism

Lecture 09 - Charges \& Currents in Magnetic Fields
Y\&F Chapter 27, Sec. 1 - 8

- What Produces Magnetic Field?
- Properties of Magnetic versus Electric Fields
- Force on a Charge Moving through Magnetic Field
- Magnetic Field Lines
- A Charged Particle Circulating in a Magnetic Field Cyclotron Frequency
- The Cyclotron, the Mass Spectrometer, the Earth's Field
- Crossed Electric and Magnetic Fields
- The e/m Ratio for Electrons
- Magnetic Force on a Current-Carrying Wire
- Torque on a Current Loop: the Motor Effect
- The Magnetic Dipole Moment
- Summary


## Magnetic Field

Electrostatic \& gravitational forces act through vector fields $\overrightarrow{\mathbf{E}}$ and $\overrightarrow{\mathbf{g}}$ Now: Magnetic force (quite different): $\overrightarrow{\mathbf{B}} \equiv$ magnetic field
Force law first: How are charges \& currents affected by a given B field? Then: How to create B field (next lecture)

- Currents in loops of wire
- Intrinsic spins of $e^{-}, p^{+} \rightarrow$ currents $\rightarrow$ magnetic dipole moment
- Spins can align permanently to form natural magnets

Unification: Maxwell's equations, electromagnetic waves

## "Permanent Magnets": North- and South- seeking poles

Natural magnets known since antiquity Earth's magnetic field, compass
Ferromagnetic materials make a magnet when cooled in B field (Fe, Ni, Co) -

- Spins align into"domains" (magnets)

Other materials (plastic, copper, wood,...) slightly or not affected (para- and dia- magnetism)


## Electric Field versus Magnetic Field

- Electric force acts at a distance through electric field.
- Vector field, E.
- Source: electric charge.
- Positive charge (+) and negative charge (-).
- Opposite charges attract
- Like charges repel.
- Electric field lines visualizing the direction and magnitude of E.

- Magnetic force acts at a distance through magnetic field.
- Vector field, B
- Source: moving electric charge (current, even in substances such as permanent magnets).
- North pole (N) and south pole (S)
- Opposite poles attract
- Like poles repel.
- Magnetic field lines visualizing the direction and magnitude of $\underline{B}$.

(b)

(c)



## Differences between magnetic \& electrostatic field

Test charge and electric field

$$
\vec{E}=\frac{\overrightarrow{F_{E}}}{\mathbf{q}}
$$

Single electric poles exist

Test monopole and magnetic field?


Magnetic poles are always found in pairs. A single magnetic pole has never been found.

Cut up a bar magnet $\square$ small, complete magnets


There is no magnetic monopole... ...dipoles are the basic units
$\begin{aligned} & \text { Electrostatic } \\ & \text { Gauss Law }\end{aligned} \int_{\mathrm{S}} \overrightarrow{\mathrm{E}} \circ \mathrm{d} \overrightarrow{\mathrm{A}}=\mathrm{q}_{\mathrm{enc}} / \varepsilon_{0} \quad \begin{aligned} & \text { Magnetic } \\ & \text { Gauss Law }\end{aligned} \int_{\mathrm{S}} \overrightarrow{\mathrm{B}} \circ \mathrm{d} \overrightarrow{\mathrm{A}}=0$
Magnetic flux through each and every Gaussian surface $=0$
Magnetic field exerts force on moving charges (current) only

## Magnetic Force on a Charged Particle

Define $\underline{B}$ by the magnetic force $E_{B}$ it exerts on a charged particle moving with a velocity $\mathbf{v}$
$\mathrm{F}_{\mathrm{B}}=\mathrm{qvB} \sin (\phi)$
$\overrightarrow{\mathrm{F}}_{\mathrm{B}}=\mathbf{q} \overrightarrow{\mathbf{v}} \times \overrightarrow{\mathbf{B}}$
"LORENTZ FORCE",

Units: $\quad[1$ Tesla $] \equiv \frac{[\text { Newtons }]}{[\text { Coulomb] }[\mathrm{m} / \mathrm{s}]}$
1 "GAUSS" = $10^{-4}$ Tesla
See Lecture 01 for cross product definitions and examples

- $F_{B}$ is proportional to speed $v$ as well as charge $q$ and field $B$.
- $F_{B}$ is geometrically complex - depends on cross product
o $F=0$ if $\underline{v}$ is parallel to $\underline{B}$.
$\circ F$ is otherwise normal to plane of both $\underline{v}$ and $B$.
o $F_{B}$ reverses sign for opposite sign of charge
o Source is also qv [current $x$ length].
- Electric force can do work on a charged particle...
- BUT magnetic force cannot do work on moving particles since $\underline{E}_{B} \cdot \underline{v}=0$.


## The Vector Product $\overrightarrow{\mathbf{c}}=\overrightarrow{\mathbf{a}} \times \overrightarrow{\mathbf{b}}$ in terms of Vector Components

$\vec{a}=a_{x} \hat{i}+a_{y} \hat{j}+a_{z} \hat{k}, \vec{b}=b_{x} \hat{i}+b_{y} \hat{j}+b_{z} \hat{k}, \vec{c}=c_{x} \hat{i}+c_{y} \hat{j}+c_{z} \hat{k}$
The vector components of vector $\vec{c}$ are given by the equations:
$c_{x}=a_{y} b_{z}-a_{z} b_{y}, \quad \mathrm{c}_{y}=a_{z} b_{x}-a_{x} b_{z}, \quad \mathrm{c}_{z}=a_{x} b_{y}-a_{y} b_{x}$

Note: Those familiar with the use of determinants can use the expression

$$
\vec{a} \times \vec{b}=\left|\begin{array}{ccc}
\hat{i} & \widehat{j} & \vec{k} \\
a_{x} & a_{y} & a_{z} \\
b_{x} & b_{y} & b_{z}
\end{array}\right|
$$

Note: The order of the two vectors in the cross product is important
$\vec{b} \times \vec{a}=-(\vec{a} \times \vec{b})$

Magnetic Field Lines : In analogy with the electric field lines we introduce the concept of magnetic field lines which help visualize the magnetic field vector $\vec{B}$ without using equations.
The relation between the magnetic field lines and $\vec{B}$ are:

1. At any point $P$ the magnetic field vector $\vec{B}$ is tangent to the magnetic field lines

2. The magnitude of the magnetic field vector $\overrightarrow{\mathbf{B}}$ is proportional to the density of the magnetic field lines

$$
B_{P}>B_{Q}
$$



## Magnetic force geometry examples:

A simple geometry

- charge +q
- $\underline{B}$ along $-z$ direction
- $\underline{v}$ along $+x$ direction

| CONVENTION |  |
| :---: | :---: |
| $\bullet$ | $X$ |
| Head | Tail |




$\begin{aligned} \therefore \vec{F}_{m} & =q \vec{v} \times \vec{B} \\ \rightarrow \quad & =q \mathbf{v}_{0} B_{0} \hat{j}\end{aligned}$


More simple examples


- charge +q
- $\underline{E}_{m}$ is out of the page
- $\mathbb{E}_{\underline{m}} \underline{L}=q \underline{V}_{o} \underline{B}_{0}$



## A numerical example

- Electron beam moving in plane of sketch
- $\mathrm{v}=10^{7} \mathrm{~m} / \mathrm{s}$ along +x
- $B=10^{-3} \mathrm{~T}$. out of page along $+y$
a) Find the force on the electron :


$$
\begin{aligned}
& \vec{F}_{m}=-e \vec{v} \times \vec{B}=-1.6 \times 10^{-19} \times 10^{7} \times 10^{-3} \times \sin \left(90^{\circ}\right) \hat{k} \\
& \vec{F}_{m}=-1.6 \times 10^{-15} \hat{k}
\end{aligned}
$$

Negative sign means force is opposite to result of using the RH rule
b) Acceleration of electron :

$$
\overrightarrow{\mathbf{a}}=\frac{\vec{F}_{\mathrm{m}}}{\mathrm{~m}_{\mathrm{e}}}=\frac{-1.6 \times 10^{-15} \mathrm{~N}}{9 \times 10^{-31} \mathrm{Kg}} \hat{\mathbf{k}}=-1.76 \times 10^{+15} \hat{\mathrm{k}}\left[\mathrm{~m} / \mathrm{s}^{2}\right]
$$

Direction is the same as that of the force

## Direction of Magnetic Force

9-2: The figures shows five situations in which a positively charged particle with velocity $\underline{v}$ travels through a uniform magnetic field $\underline{B}$. For which situation is the direction of the magnetic force along the $+x$ axis ?
Hint: Use Right Hand Rule.


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## Magnetic Units and Field Line Examples

- SI unit of magnetic field: tesla ( T )
$-1 \mathrm{~T}=1 \mathrm{~N} /[\mathrm{Cm} / \mathrm{s}]=1 \mathrm{~N} /[\mathrm{Am}]=10^{4}$ gauss
- Magnetic field lines - interpret similarly to E
- The tangent to a magnetic field line at any point gives the direction of $B$ at that point;
- The spacing of the lines represents the magnitude of $B-$ the magnetic field is stronger where the lines are closer together, and conversely.


| At surface of neutron star | $10^{8} \mathrm{~T}$ |
| :--- | :--- |
| Near big electromagnet | 1.5 T |
| Inside sunspot | $10^{-1} \mathrm{~T}$ |
| Near small bar magnet | $10^{-2} \mathrm{~T}$ |
| At Earth's surface | $10^{-4} \mathrm{~T}$ |
| In interstellar space | $10^{-10} \mathrm{~T}$ |



## Charged Particles Circle at Constant Speed in a Uniform

 Magnetic Field: Cyclotron Frequency- Uniform B in z direction, $\underline{v}$ in $x-y$ plane tangent to path
- $\underline{F}_{B}$ is normal to both $\underline{v} \& B$... so Power $=\underline{F} \cdot \underline{v}=0$.
- Magnetic force can not change particle's speed or KE
- $\underline{F}_{B}$ is a centripetal force, motion is UCM
- A charged particle moving in a plane perpendicular to a B field circles in the plane with constant speed v

| Set | $F_{B}=q v B=\frac{\mathbf{m} v^{2}}{}$ |
| :---: | :---: |
| Radius of the path | $r=\frac{m v}{q B}$ |
| Period | $\tau_{c}=\frac{2 \pi r}{v}=\frac{2 \pi m}{q B}$ |
| Cyclotron angular frequency | $\omega_{c} \equiv \frac{2 \pi}{\tau_{c}}=\frac{q B}{m}$ |

- $t$ and $\omega$ do not depend on velocity.
- Fast particles move in large circles and slow ones in small circles
- All particles with the same charge-to-mass ratio have the same period.

If $\underline{v}$ not normal to $\underline{B}$ particle spirals around $B$
$\overrightarrow{\mathbf{v}}_{\perp}$ causes circling
$\mathbf{V}_{\text {para }}$ is cons tan $t$

- The rotation direction for a positive and negative particles is opposite.


## Cyclotron particle accelerator

Early nuclear physics research, Biomedical applications


- Inject charged particles in center
- Charged "Dees" reverse polarity as particles cross gap to accelerate them.
- Particles spiral out in magnetic field as they gain KE and are detected

$$
r=\frac{\mathbf{m v}}{\mathbf{q B}}
$$

- Frequency of polarity reversal needed does not depend on speed or radius of path - it can be constant!

$$
\omega_{c}=\frac{2 \pi}{\tau_{c}}=\frac{q B}{m}
$$

## Earth's field shields us from the Solar Wind and produces the Aurora

Earth's magnetic field deflects charged solar wind particles (cyclotron effect) protecting the Earth and making life possible (magnetosphere).


## Another cyclotron effect device: Mass spectrometer

 Separates particles with different charge/mass ratios

## Circulating Charged Particle

9-3: The figures show circular paths of two particles having the same speed in a uniform magnetic field B , which is directed into the page. One particle is a proton; the other is an electron (much less massive). Which figure is physically reasonable?


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## Charged particle in both $\underline{E}$ and $\underline{B}$ fields

Add the electrostatic and magnetic forces

$$
\overrightarrow{\mathbf{F}}_{\text {tot }}=\mathbf{q} \overrightarrow{\mathbf{E}}+\mathbf{q} \overrightarrow{\mathbf{v}} \times \overrightarrow{\mathbf{B}}
$$

$$
\begin{aligned}
& \vec{F}_{\text {tot }}=m \vec{a} \\
& \text { also }
\end{aligned}
$$

Does $F=$ ma change if you observe this while moving at constant velocity v'?
Example: Crossed E and B fields



- B out of paper
- $\underline{E}$ up and normal to $\underline{B}$
-     + charge
- $\underline{v}$ normal to both $\underline{E} \& \underline{B}$

$$
\begin{aligned}
& \vec{F}_{e}=q \vec{E} \quad \text { (up) } \\
& F_{m}=q \overrightarrow{\mathrm{v}} \times \vec{B} \quad \text { (down) }
\end{aligned}
$$



FBD of $q$

Equilibrium when...

$$
\vec{F}_{\text {tot }}=0 \quad \Longrightarrow \quad q E=q v B \quad v=|E| /|B|
$$

OPPOSED
FORCES


- Independent of charge q
- Charges move through fields un-deflected
- Use to select particles with a particular velocity

Measuring e/m ratio for the electron (J. J. Thompson, 1897)


First: . Apply E field only in -y direction
. $\mathrm{y}=$ deflection of beam from center of screen (position for $\mathrm{E}=0$ )

$$
\begin{aligned}
F_{y}=q E & =m a_{y} \\
q= & -e \\
& \text { measure } y=\frac{1}{2} a_{y} t^{2} \\
& \text { notethat } t=L / v_{x}
\end{aligned}
$$

$$
\frac{e}{m}=\frac{a_{y}}{E}
$$

Next: . Find $\mathrm{a}_{\mathrm{y}}$ by measuring y and flight time $\mathrm{t}=\mathrm{L} / \mathrm{v}_{\mathrm{x}}$ (constant)
Use crossed $B$ and $E$ fields to measure $v_{x}$
Let $B$ point into the page, perpendicular to both $E$ and $V x$
$F_{M}$ points along $-y$, opposite to $F E$ (negative charge)

$$
\vec{F}_{M}=-\mathrm{e} \overrightarrow{\mathbf{v}} \times \overrightarrow{\mathbf{B}} \quad \overrightarrow{\mathrm{F}}_{\mathrm{E}}=-\mathrm{e} \overrightarrow{\mathbf{E}}
$$

$$
\text { at equilibrium } \quad v_{x}=E / B
$$

Adjust $B$ until beam deflection $=0\left(F_{E}\right.$ cancels $\left.F_{M}\right)$ to find $v_{x}$ and time $\dagger$

$$
t=\frac{B L}{E} \Rightarrow a_{y}=\frac{2 y}{t^{2}}=2 y \frac{E^{2}}{B^{2} L^{2}}
$$

$$
\frac{e}{m}=2 y \frac{E}{B^{2} L^{2}}=1.76 \times 10^{11} \mathrm{C} / \mathrm{kg}
$$

## Force due to crossed E and B fields

9-4: The figure shows four directions for the velocity vector $\underline{v}$ of a positively charged particle moving through a uniform electric field $\underline{E}$ (into the page) and a uniform magnetic field $\underline{B}$ (pointing to the right). The speed is $E / B$.
Which direction of velocity produces the greatest magnitude of the net force?


## Force on a straight wire carrying current in a $\underline{B}$ field

Free electrons (negative):

- Drift velocity $\mathrm{v}_{\mathrm{d}}$ is opposite to current (along wire)
- Lorentz force on an electron $=-\mathrm{ev}_{d} \times \underline{\mathrm{B}} \quad$ (normal to wire) Motor effect: wire is pushed or pulled by the charges

$$
\begin{aligned}
& \overrightarrow{\mathbf{v}}_{\mathbf{d}}=-\left|\mathbf{v}_{\mathbf{d}}\right| \hat{\mathbf{n}} \quad \hat{\mathbf{n}} \equiv \text { unit vectoralongcurrent } \\
& \mathbf{d} \overrightarrow{\mathrm{F}}_{\mathbf{m}} \equiv \text { forcedue to charges in length } \mathbf{d} \overrightarrow{\mathbf{x}} \\
& \mathbf{d} \vec{F}_{\mathbf{m}}=\mathbf{d q} \mathbf{q}_{\times} \times \overrightarrow{\mathbf{v}} \quad \begin{array}{c}
\text { CONVEN } \\
\text { OUT }
\end{array} \\
& \begin{array}{l}
\text { dq is the charge moving in wire } \\
\text { whose length } d x=-v_{d} d t
\end{array}
\end{aligned}
$$



Recall: $d q=-i d t$, Note: $d \underline{x}$ is in direction of current (+ charges)

$$
d \vec{F}_{m}=-i\left(\vec{v}_{d} d t\right) \times \vec{B}=i d \vec{x} \times \vec{B}
$$

Integrate along the whole length $L$ of the wire (assume $\underline{B}$ is constant )

$$
\overrightarrow{\mathrm{L}}=\text { lengthof wire, parallelto current }
$$



$$
\vec{F}_{m}=i \vec{L} \times \vec{B}
$$

## Motor effect on a wire: which direction is the pull?

## $\vec{F}_{\mathrm{m}}=\mathrm{i} \overrightarrow{\mathbf{L}} \times \vec{B}$

i $\overrightarrow{\mathbf{L}}$ replaces $\mathbf{q} \overrightarrow{\mathbf{V}}$


A current-carrying loop experiences A torque (but zero net force)


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Example: A 3.0 A. current flows along $+x$ in a wire 1.0 m long aligned parallel to the $x$-axis. Find the magnitude and direction of the magnetic force on the wire, assuming that the magnetic field is uniform and given by:

$$
\begin{gathered}
\overrightarrow{\mathbf{B}}=\mathbf{3 . 0} \hat{\mathbf{i}}+5.0 \hat{\mathbf{j}} \text { [Tesla] } \quad i \overrightarrow{\mathbf{L}}=3.0 \times 1.0 \hat{\mathbf{i}}[\mathrm{~m}] \\
\overrightarrow{\mathbf{F}}_{\mathrm{m}}= \\
\mathbf{i} \overrightarrow{\mathbf{L}} \times \overrightarrow{\mathbf{B}}=\mathbf{3 . 0 \times 1 . 0} \hat{\mathbf{i}} \times(\mathbf{3 . 0} \hat{\mathbf{i}}+5.0 \hat{\mathbf{j}}) \\
\hat{\mathbf{i}} \times \hat{\mathbf{i}}=0 \quad \hat{\mathbf{i}} \times \hat{\mathbf{j}}=\hat{\mathbf{k}}
\end{gathered}
$$

$$
\therefore \vec{F}_{\mathbf{m}}=15.0 \hat{\mathbf{k}}[\mathrm{~N}]
$$



Only the component of B perpendicular to the current-length contributes to the force

Example: Same as above, but now the field has a z-component

$$
\begin{aligned}
& \vec{B}=3.0 \hat{i}+5.0 \hat{\mathbf{j}}+4.0 \hat{\mathbf{k}}{ }_{\text {[Tesla] } \quad i \vec{L}=3.0 \times 1.0 \hat{i}[\mathrm{~m}]} \\
& \overrightarrow{\mathbf{F}}_{\mathrm{m}}=\mathrm{i} \overrightarrow{\mathrm{~L}} \times \overrightarrow{\mathbf{B}}=3.0 \times 1.0 \hat{\mathbf{i}} \times(3.0 \hat{\mathbf{i}}+5.0 \hat{\mathbf{j}}+4.0 \hat{\mathbf{k}})
\end{aligned}
$$

$\hat{\mathbf{i}} \mathbf{x} \hat{\mathbf{i}}=\mathbf{0} \quad \hat{\mathbf{i}} \times \hat{\mathbf{j}}=\hat{\mathbf{k}} \quad \hat{\mathbf{i}} \times \hat{\mathbf{k}}=-\hat{\mathbf{j}}$
$\therefore \vec{F}_{\mathrm{m}}=15.0 \hat{\mathbf{k}}-12.0 \hat{\mathrm{j}}[\mathrm{N}]$

Whatever the direction of the field $B$, there can never be a force component along the current-length direction

Torque on a current-carrying wire loop in a magnetic field



- Upper sketch: x-axis toward viewer
- Loop can rotate about $x$-axis $B$ field is along $z$ axis
- Apply $\overrightarrow{\boldsymbol{F}}_{\mathbf{m}}=\mathbf{i} \overrightarrow{\mathbf{L}} \times \overrightarrow{\mathbf{B}}$ to each side of the loop
- $\left|F_{1}\right|=\left|F_{3}\right|=$ iaB: Forces cancel again but net torque is not zero!
- Moment arms for $F_{1} \& F_{3}$ equal b. $\sin (q) / 2$
- Force $F_{1}$ produces CW torque equal to

$$
t_{1}=\text { i.a.B.b. } \sin (q) / 2
$$

- Same for $F_{3}$
- Torque vector is down into paper along $-x$ rotation axis

Net torque $\tau=\operatorname{iabBsin}(\theta)$

- Lower sketch:
- $\left|F_{2}\right|=\left|F_{4}\right|=$ i.b. $B \cos (\theta)$ : Forces cancel, same line of action $\rightarrow$ zero torque


## Torque on a current-carrying loop: Motor Effect


torque $\tau=\operatorname{iabBsin}(\theta) \longmapsto \vec{\tau}=\mathbf{i} \mathbf{A} \hat{\mathbf{n}} \times \overrightarrow{\mathbf{B}} \quad \mathbf{A} \equiv$ area of loop $=\mathbf{a b}$

Maximum torque: $\tau_{\max }=\mathbf{i} \mathbf{A B} \quad \hat{\mathbf{n}} \perp \overrightarrow{\mathbf{B}}$
Sinusoidal variation: $\tau(\theta)=\tau_{\text {max }} \sin (\theta)$ Stable when $\hat{\mathbf{n}}$ parallels $\overrightarrow{\mathbf{B}}$
Restoring torque $\rightarrow$ oscillations

Field PRODUCED by a current loop is a magnetic dipole field

## ELECTRIC MOTOR:

- Reverse current I when torque $\tau$ changes $\operatorname{sign}(\theta=0)$
- Use mechanical "commutator"

- Multiple turns of wire increase torque
- N turns assumed in flat, planar coil
- Real motors use multiple coils (smooth torque)

$$
\begin{aligned}
& \vec{\tau}=\operatorname{NiA} \hat{\mathbf{n}} \times \vec{B} \quad \mathbf{A} \equiv \text { area of loop }=\mathrm{ab} \\
& \tau=\operatorname{NiA} \mathrm{B} \sin (\theta)
\end{aligned}
$$

## MOVING COIL GALVANOMETER

- Basis of most $19^{\text {th }} \& 20^{\text {th }}$ century analog instruments: voltmeter, ohmmeter, ammeter, speedometer, gas gauge, ....
- Coil spring calibrated to balance torque at proper mark on scale


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## Current loops are basic magnetic dipoles

Represent loop as a vector

## Magneticdipole moment $\equiv \vec{\mu} \equiv \mathbf{N i} \mathbf{A} \hat{\mathbf{n}}$



Dimensions $[\mu]=$ ampere- $\mathrm{m}^{2}=\frac{\text { Newton }-\mathrm{m}}{\text { Tesla }}=\frac{\text { Joule }}{\text { Tesla }}$ $\mathrm{N}=$ numberof turnsin theloop

$$
\begin{aligned}
& \vec{\tau}=\vec{\mu} \times \overrightarrow{\mathbf{B}} \quad \mathbf{A} \equiv \text { area of loop }=\mathbf{a b} \\
& \tau=\mu \operatorname{Bin}(\theta)
\end{aligned}
$$



$$
\mathrm{U}_{\mathrm{M}}=\int \tau \mathrm{d} \theta=\int \mu \mathrm{B} \sin (\theta) \mathrm{d} \theta
$$

Magnetic dipole moment $\underline{\mu}$ measures


SAMPLE VALUES [J / T]

$\mathrm{m} \sim 8.0 \times 10^{22} \mathrm{~J} / \mathrm{T}$
torque
$\vec{\tau}_{\mathbf{e}}=\overrightarrow{\mathbf{p}} \times \overrightarrow{\mathbf{E}}$
$\vec{\tau}_{\mathrm{m}}=\vec{\mu} \times \overrightarrow{\mathbf{B}}$
POTENTIAL ENERGY
$\mathbf{U}_{\mathbf{e}}=-\overrightarrow{\mathbf{p}} \cdot \overrightarrow{\mathbf{E}}$
$\mathbf{U}_{\mathrm{m}}=-\vec{\mu} \circ \overrightarrow{\mathbf{B}}$

## Torque and potential energy of a dipole

9-5: In which configuration is the torque on the dipole at it's maximum value?

$$
\vec{\tau}_{\mathbf{m}}=\vec{\mu} \times \overrightarrow{\mathbf{B}}
$$

9-6: In which configuration is the potential energy of the dipole at it's smallest value?

$$
\mathrm{U}_{\mathrm{m}}=-\vec{\mu} \circ \overrightarrow{\mathbf{B}}
$$



## Summary: Lecture 9 Chapter 28 - Magnetic Fields

## chapter 27 sUMMARY

Magnetic forces: Magnetic interactions are fundamentally interactions between moving charged particles. These interactions are described by the vector magnetic field, denoted by $\overrightarrow{\boldsymbol{B}}$. A particle with charge $q$ moving with velocity $\overrightarrow{\boldsymbol{v}}$ in a magnetic field $\overrightarrow{\boldsymbol{B}}$ experiences a force $\overrightarrow{\boldsymbol{F}}$ that is perpendicular to both $\overrightarrow{\boldsymbol{v}}$ and $\overrightarrow{\boldsymbol{B}}$. The SI unit of magnetic field is the tesla ( $1 \mathrm{~T}=1 \mathrm{~N} / \mathrm{A} \cdot \mathrm{m}$ ). (See Example 27.1.)
$\vec{F}=q \vec{v} \times \vec{B}$
(27.2)


Magnetic field lines and flux: A magnetic field can be represented graphically by magnetic field lines. At each point a magnetic field line is tangent to the direction of $\overrightarrow{\boldsymbol{B}}$ at that point. Where field lines are close together the field magnitude is large, and vice versa. Magnetic flux $\Phi_{B}$ through an area is defined in an analogous way to electric flux. The SI unit of magnetic flux is the weber ( $1 \mathrm{~Wb}=1 \mathrm{~T} \cdot \mathrm{~m}^{2}$ ). The net magnetic flux through any closed surface is zero (Gauss's law for magnetism). As a result, magnetic field lines always close on themselves. (See Example 27.2.)
$\Phi_{B}=\int B_{\perp} d A$
$=\int B \cos \phi d A$
$=\int \overrightarrow{\boldsymbol{B}} \cdot d \overrightarrow{\boldsymbol{A}}$
$\oint \overrightarrow{\boldsymbol{B}} \cdot d \overrightarrow{\boldsymbol{A}}=0$ (closed surface)

$R=\frac{m v}{|q| B}$


Motion in a magnetic field: The magnetic force is always perpendicular to $\overrightarrow{\boldsymbol{v}}$; a particle moving under the action of a magnetic field alone moves with constant speed. In a uniform field, a particle with initial velocity perpendicular to the field moves in a circle with radius $R$ that depends on the magnetic field strength $B$ and the particle mass $m$, speed $v$, and charge $q$. (See Examples 27.3 and 27.4.)

Crossed electric and magnetic fields can be used as a velocity selector. The electric and magnetic forces exactly cancel when $v=E / B$. (See Examples 27.5 and 27.6.)

Magnetic force on a conductor: A straight segment of a conductor carrying current $I$ in a uniform magnetic field $\overrightarrow{\boldsymbol{B}}$ experiences a force $\overrightarrow{\boldsymbol{F}}$ that is perpendicular to both $\overrightarrow{\boldsymbol{B}}$ and the vector $\vec{l}$, which points in the direction of the current and has magnitude equal to the length of the segment. A similar relationship gives the force $d \overrightarrow{\boldsymbol{F}}$ on an infinitesimal current-carrying segment $d \overrightarrow{\boldsymbol{l}}$. (See Examples 27.7 and 27.8.)

$$
\begin{align*}
& \overrightarrow{\boldsymbol{F}}=\vec{I} \times \overrightarrow{\boldsymbol{B}}  \tag{27.19}\\
& d \overrightarrow{\boldsymbol{F}}=I d \vec{l} \times \overrightarrow{\boldsymbol{B}} \tag{27.20}
\end{align*}
$$

Magnetic torque: A current loop with area $A$ and current $I$ in a uniform magnetic field $\boldsymbol{B}$ experiences no net magnetic force, but does experience a magnetic torque of magnitude $\tau$. The vector torque $\overrightarrow{\boldsymbol{\tau}}$ can be expressed in terms of the magnetic moment $\overrightarrow{\boldsymbol{\mu}}=I \overrightarrow{\boldsymbol{A}}$ of the loop, as can the potential energy $U$ of a magnetic moment in a magnetic field $\boldsymbol{B}$. The magnetic moment of a loop depends only on the current and the area; it is independent of the shape of the loop. (See Examples 27.9 and 27.10.)

$$
\begin{aligned}
\tau & =I B A \sin \phi \\
\overrightarrow{\boldsymbol{\tau}} & =\overrightarrow{\boldsymbol{\mu}} \times \overrightarrow{\boldsymbol{B}} \\
U & =-\overrightarrow{\boldsymbol{\mu}} \cdot \overrightarrow{\boldsymbol{B}}=-\mu B \cos \phi
\end{aligned}
$$



